

# Spin Content of the Nucleon

or

What We've Learned from Polarized Electron Scattering

Karl J. Slifer  
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# Spin Content of the Nucleon

or

What We've Learned from Polarized Electron Scattering,



*I've*

*(in the last few months)*

Karl J. Slifer  
University of New Hampshire

# This Talk

## Burkhardt–Cottingham Sum Rule

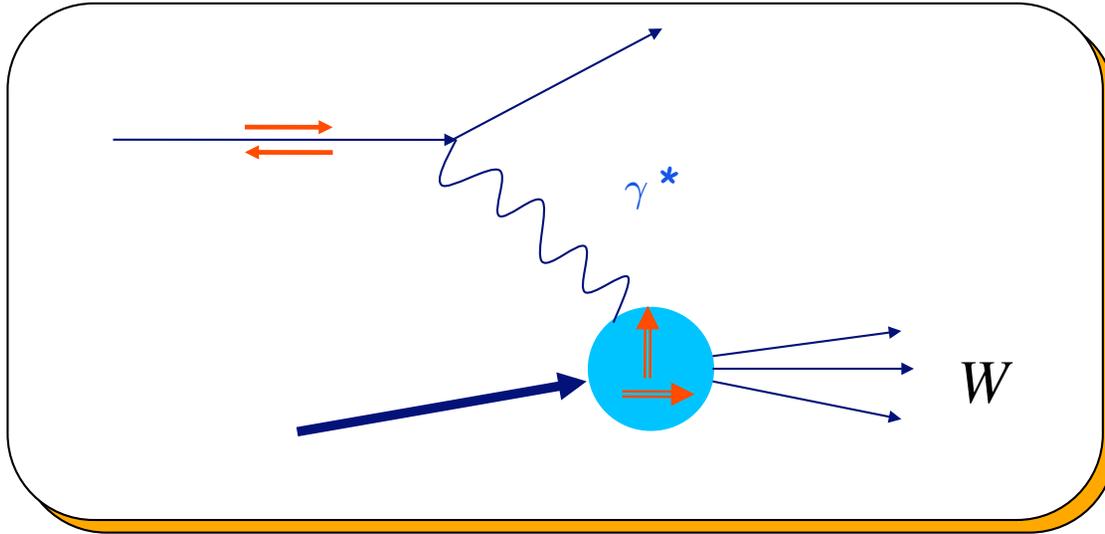
What does the JLab data tell us?  
Is it enough to make a definitive statement?

## Higher Twist Measurements at Jlab

## Target Mass Corrections

impact on the clean extraction of Higher Twist

# Inclusive Electron Scattering



When we add spin degrees of freedom to the target and beam, 2 Additional SF needed.

$$\frac{d^2\sigma}{d\Omega dE'} = \sigma_{Mott} \left[ \frac{1}{\nu} F_2(x, Q^2) + \frac{2}{M} F_1(x, Q^2) \tan^2 \frac{\theta}{2} \right] + \gamma g_1(x, Q^2) + \delta g_2(x, Q^2)$$

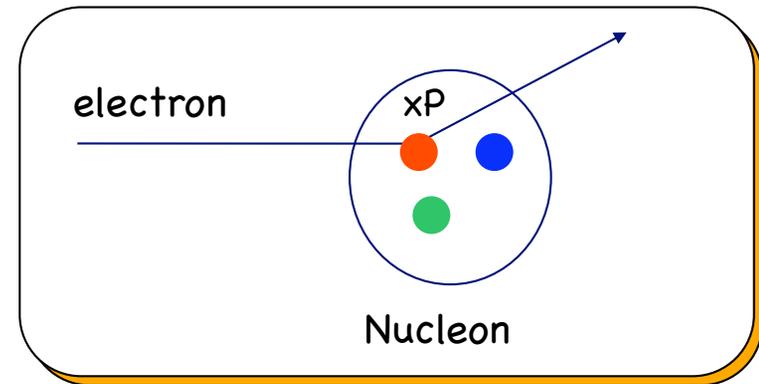
Inclusive Polarized  
Cross Section

all four SF needed for a complete description of nucleon structure

# Parton Model

## Interpretation of the Structure Functions

Impulse Approximation in DIS  
no time for interaction between partons



distributions of quark momentum and spin in the nucleon.

$$F_1(x) = \frac{1}{2} \sum e_i^2 [q_i(x) + \bar{q}_i(x)]$$

runs over all quark flavors

$$F_2(x) = 2xF_1(x)$$

$$g_1(x) = \frac{1}{2} \sum e_i^2 \Delta q_i(x)$$

$$g_2(x) = ???$$

# Burkhardt-Cottingham Sum Rule

$$\int_0^1 g_2(x, Q^2) dx = 0$$

H.Burkhardt and W.N. Cottingham  
Annals Phys. 56 (1970) 453.

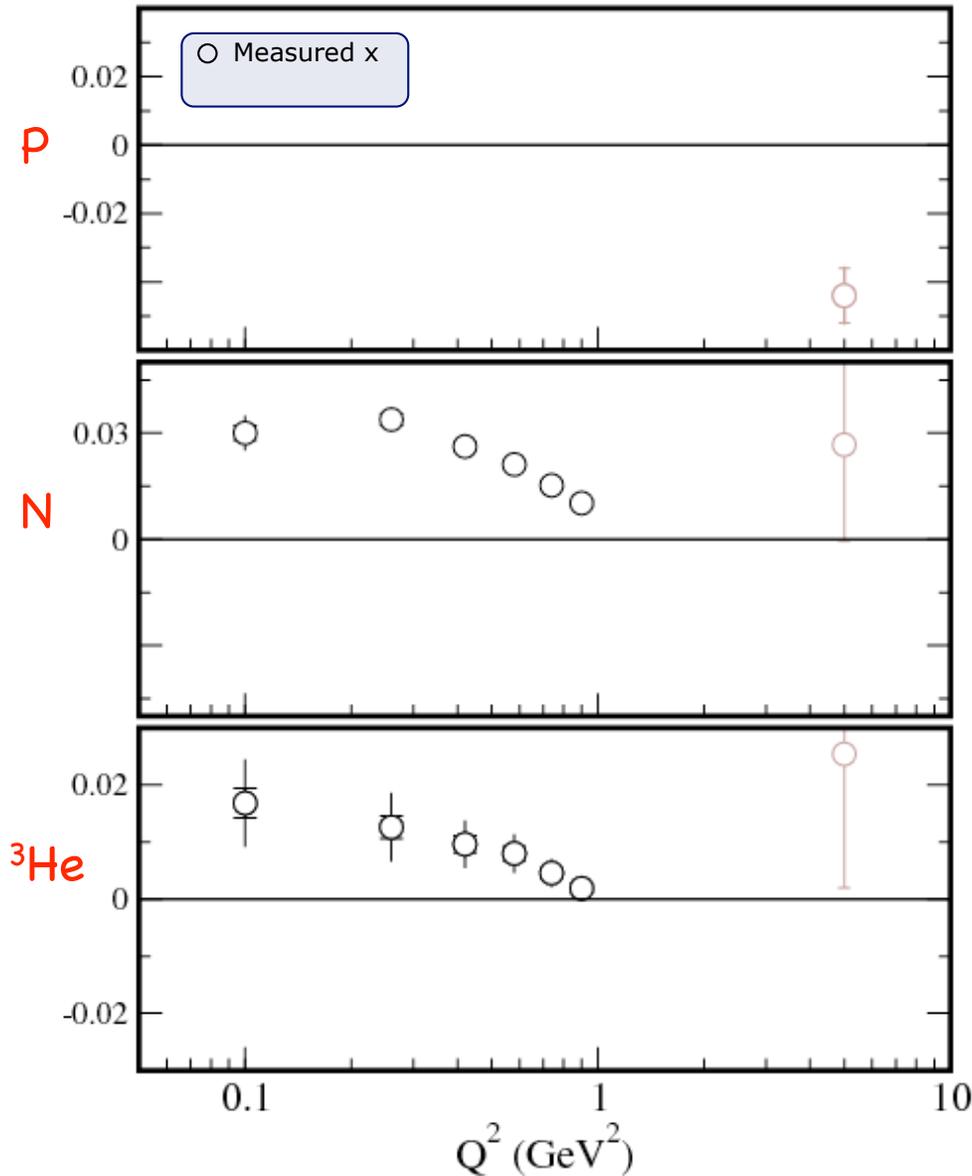
Relies on the virtual Compton scattering amplitude  $S_2$  falling to zero faster than  $1/\nu$  as  $\nu \rightarrow \infty$

Discussion of possible causes of violations

R.L. Jaffe Comm. Nucl. Part. Phys. 19, 239 (1990)

“If it holds for one  $Q^2$  it holds for all”

# BC Sum Rule



## Existing World Data on $\Gamma_2$

for Proton Neutron and  $^3\text{He}$

**BLACK** : E94010. (Hall A,  $^3\text{He}$ )

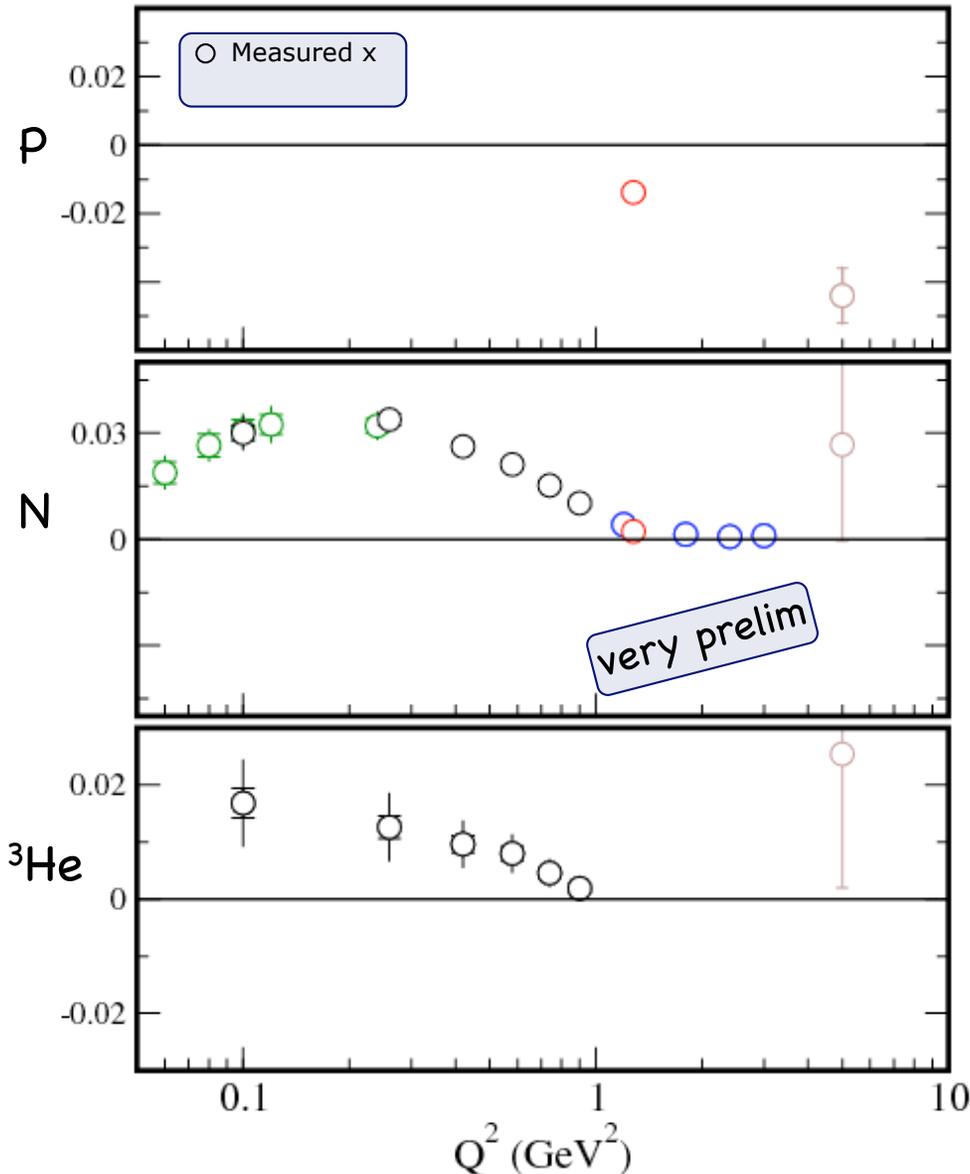
**BROWN** : E155. (SLAC NH3,  $^6\text{LiD}$ )

Note:

SLAC "Measured" =  $0.02 < x < 0.8$

JLAB "Measured"  $\approx$  Resonance Region

# BC Sum Rule



## BRAND NEW DATA!

### Very Preliminary

RED : RSS. (Hall C, NH<sub>3</sub>,ND<sub>3</sub>)

K. Slifer, O. Rondon *et al.* in preparation

BLUE: E01-012. (Hall A, <sup>3</sup>He)

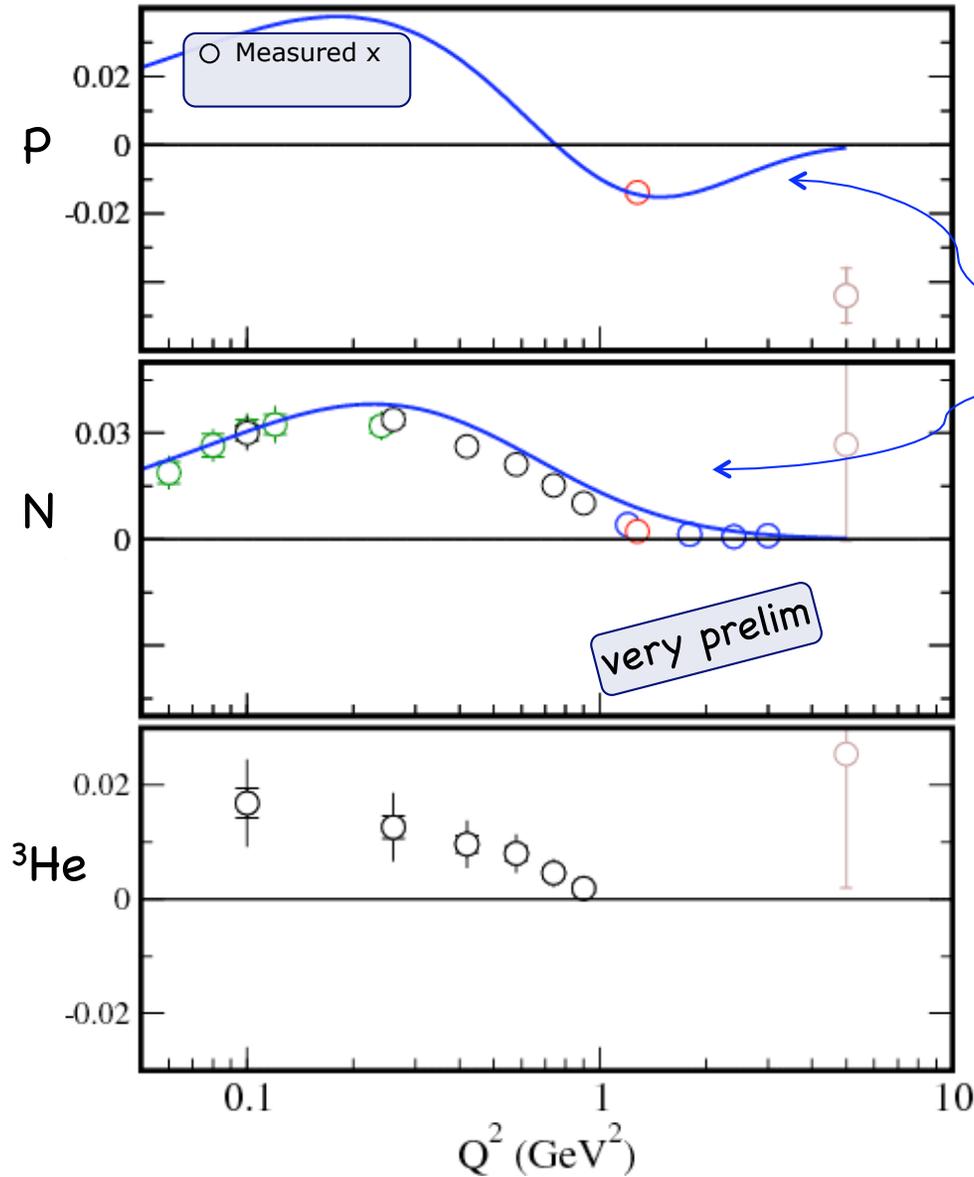
courtesy of [P. Solvignon](#)

GREEN: E97-110. (Hall A, <sup>3</sup>He)

courtesy of [V. Sulkosky](#)

Thanks also to the spokesmen of these experiments!  
RSS: Mark Jones, Oscar Rondon  
E01-012: N. Liyanage, J.P.Chen, Seonho Choi  
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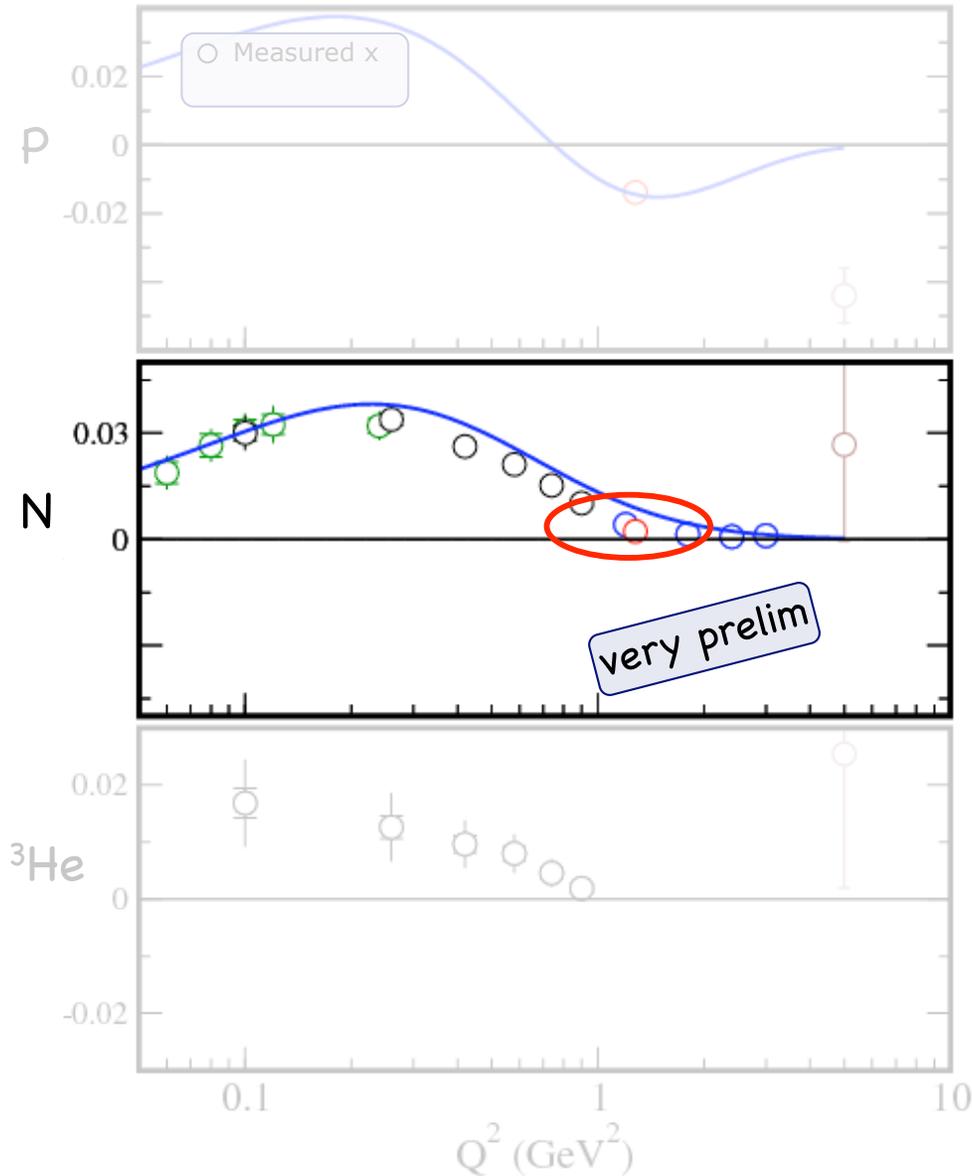
# BC Sum Rule



Good agreement with MAID model of resonance region for JLab data

very prelim

# BC Sum Rule



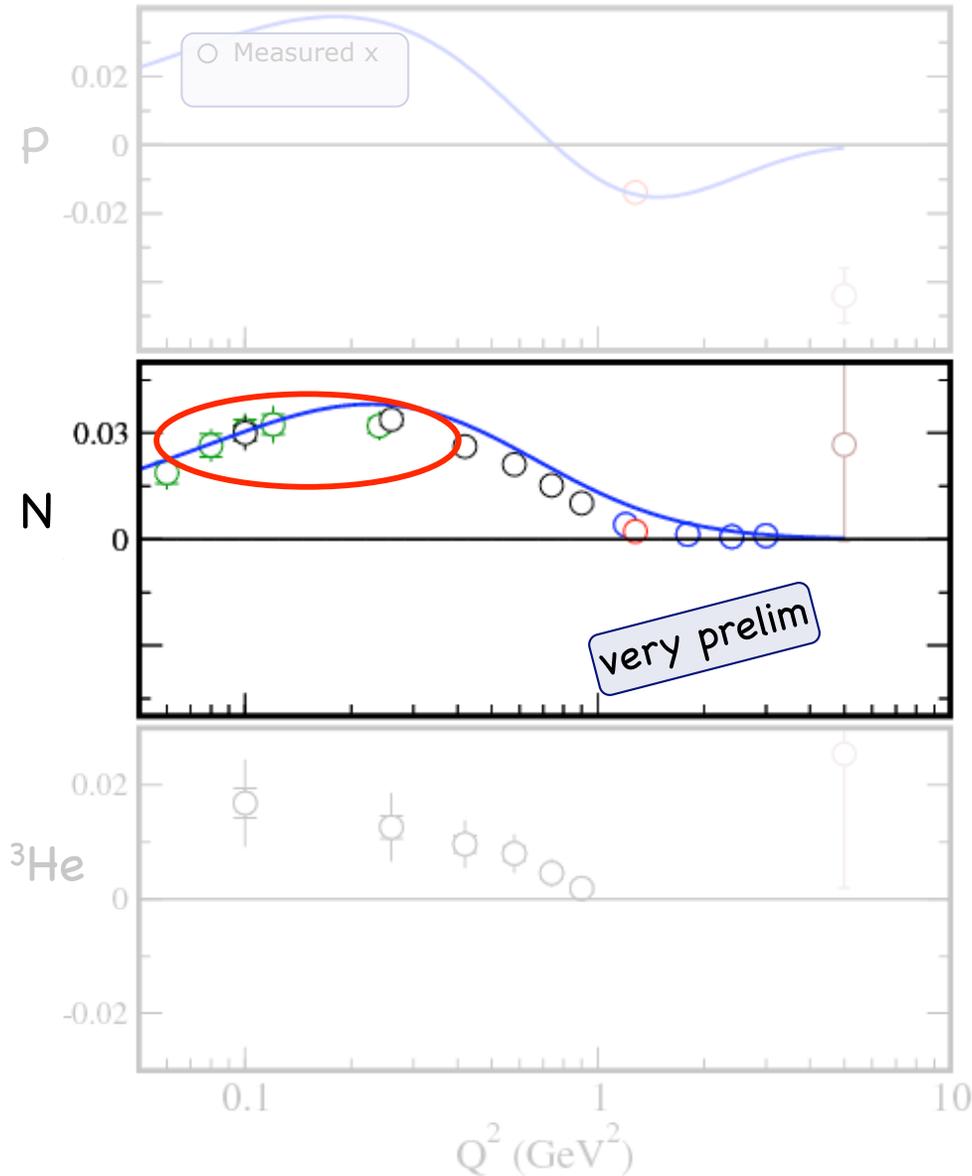
Neutron results around  $Q^2=1.3 \text{ GeV}^2$   
from 2 very different experiments:

**RSS** in Hall C: Neutron from  $\text{ND}_3$  &  $\text{NH}_3$

**E01-012** in Hall A : Neutron from  $^3\text{He}$

Excellent agreement!

# BC Sum Rule

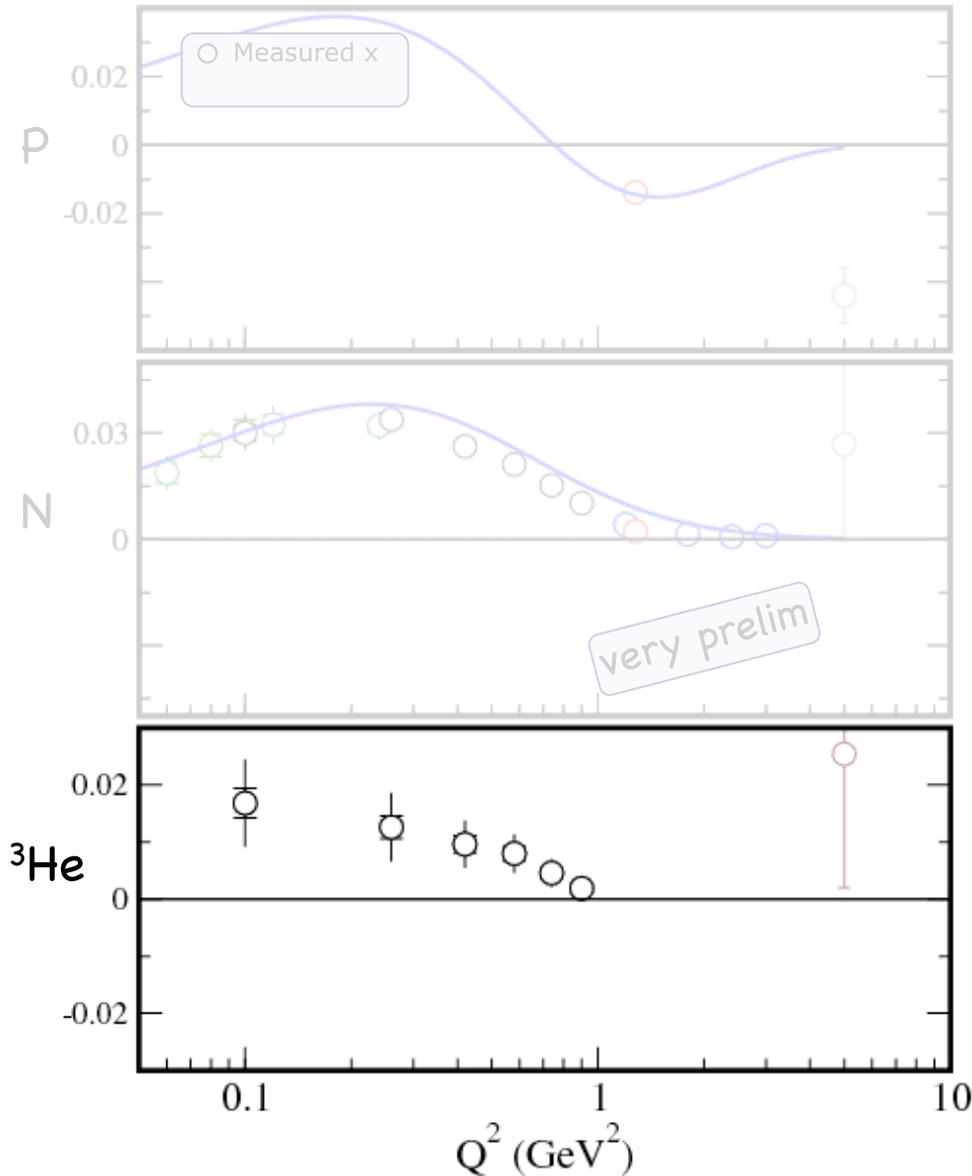


Good overlap at low  $Q^2$  of the old and new neutron data

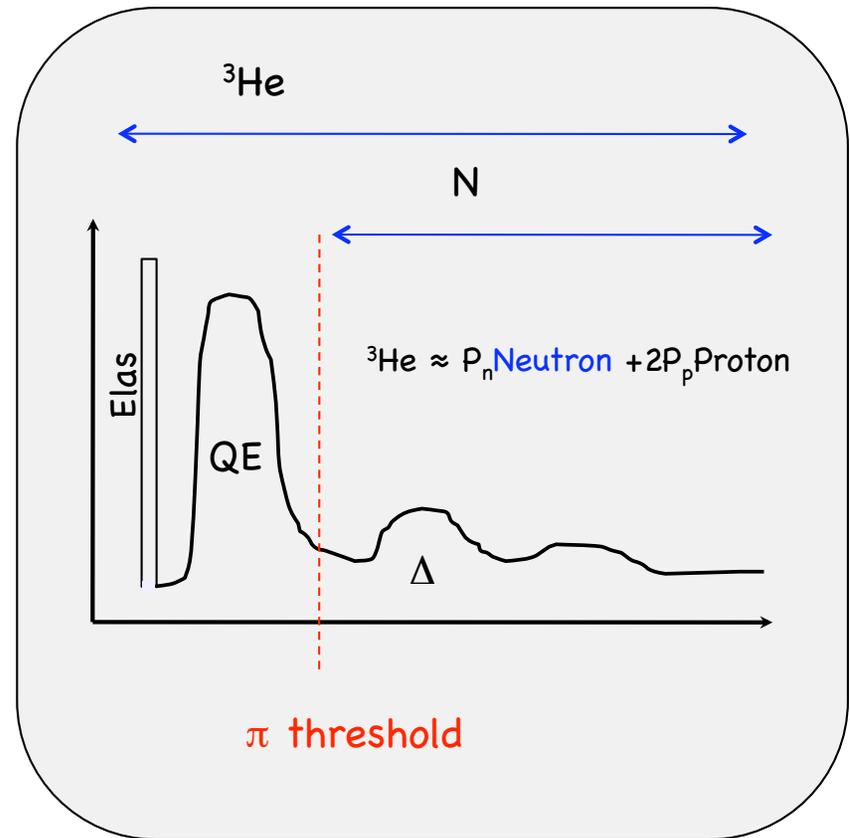
E94010 : Hall A  $^3\text{He}$  old

E97-110 Hall A  $^3\text{He}$  new

# BC Sum Rule



Difference between N and  $^3\text{He}$  Sum rules



Note:  $^3\text{He}$  requires use of nuclear elastic

# BC Sum Rule

$$\int_0^1 g_2(x, Q^2) dx = 0$$

BC = RES+DIS+ELASTIC

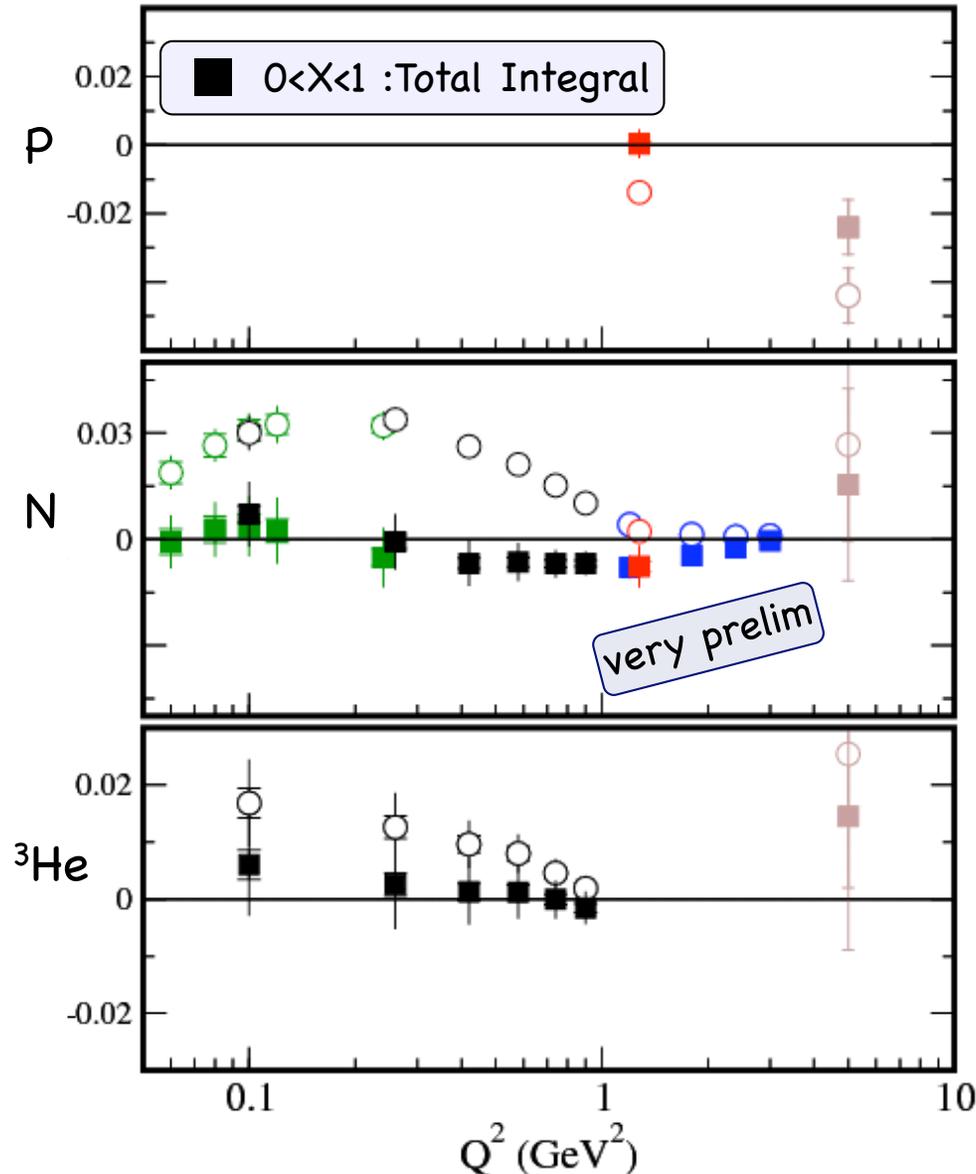
"RES": Here refers to measured x-range

"DIS": refers to unmeasured low x part of the integral. Not strictly Deep Inelastic Scattering due to low  $Q^2$

Assume Leading Twist Behaviour

Elastic: From well know FFs (<5%)

# BC Sum Rule



$$\underline{BC = RES + DIS + ELASTIC}$$

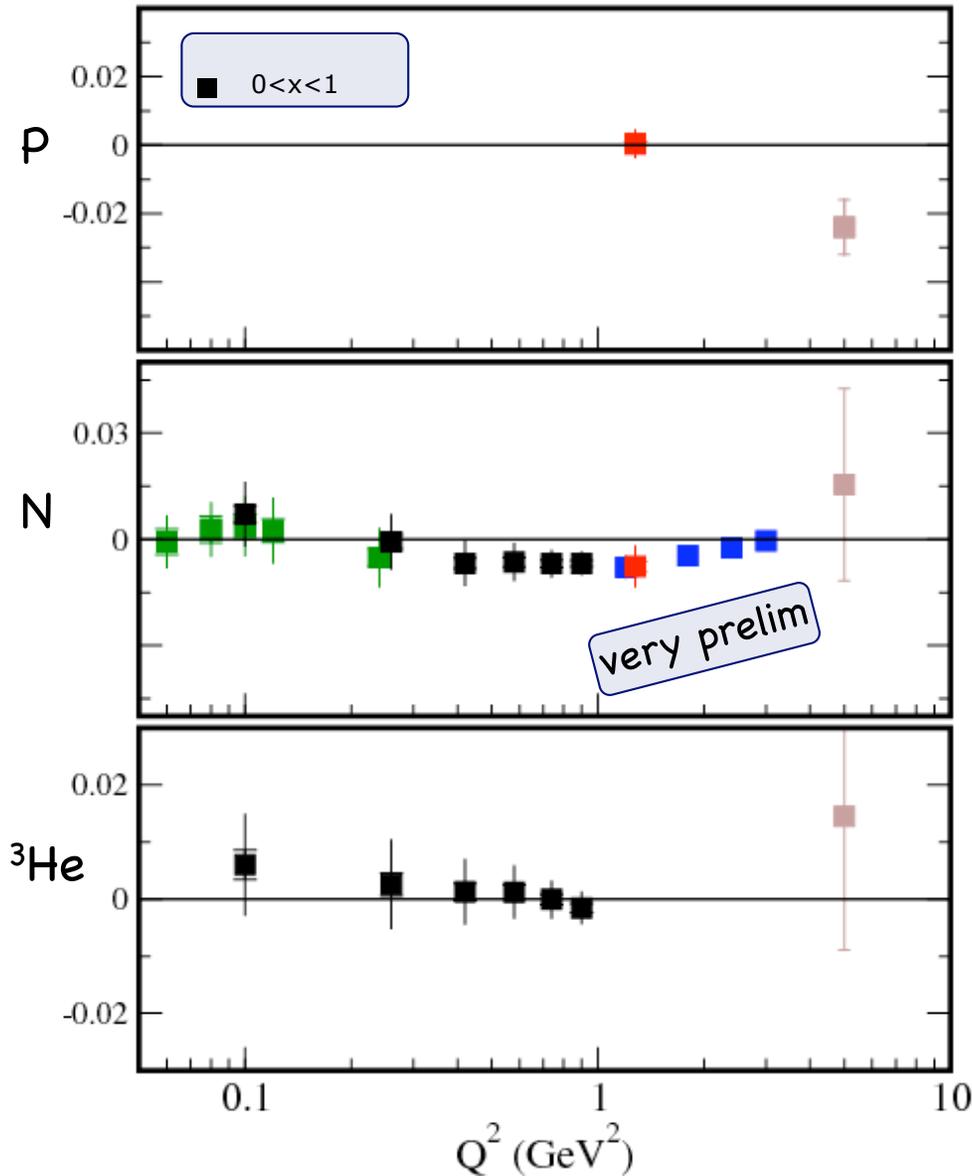
"RES": Here refers to measured x-range

"DIS": refers to unmeasured low x part of the integral. Not strictly Deep Inelastic Scattering due to low  $Q^2$

Assume Leading Twist Behaviour

Elastic: From well know FFs (<5%)

# BC Sum Rule

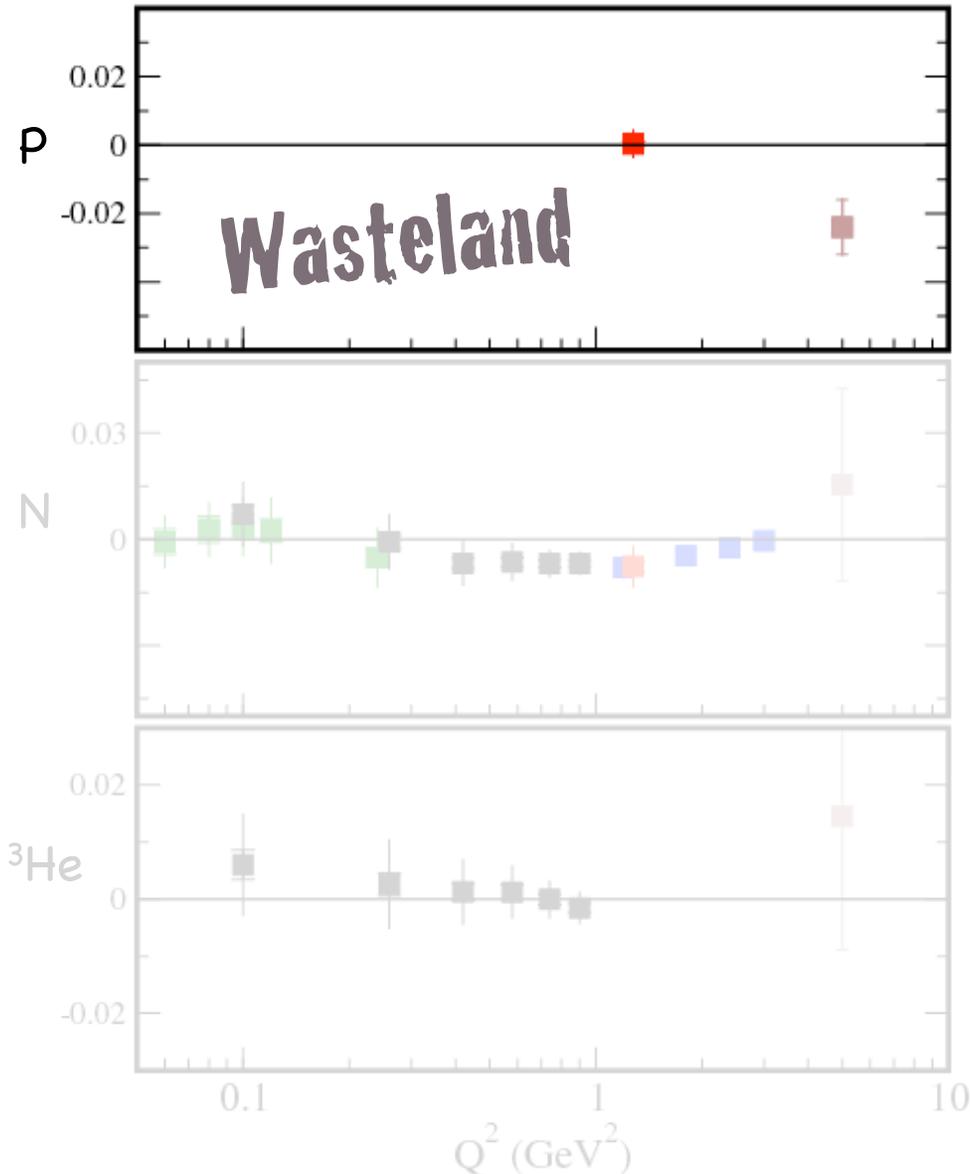


BC satisfied w/in errors for JLab Proton  
2.8 $\sigma$  violation seen in SLAC data

BC satisfied w/in errors for Neutron  
(But just barely in vicinity of  $Q^2=1$ !)

BC satisfied w/in errors for <sup>3</sup>He

# BC Sum Rule



Proton  $g_{2p}$  still relatively unknown for such a fundamental quantity.

Need more high quality data like RSS

## Upcoming Experiments

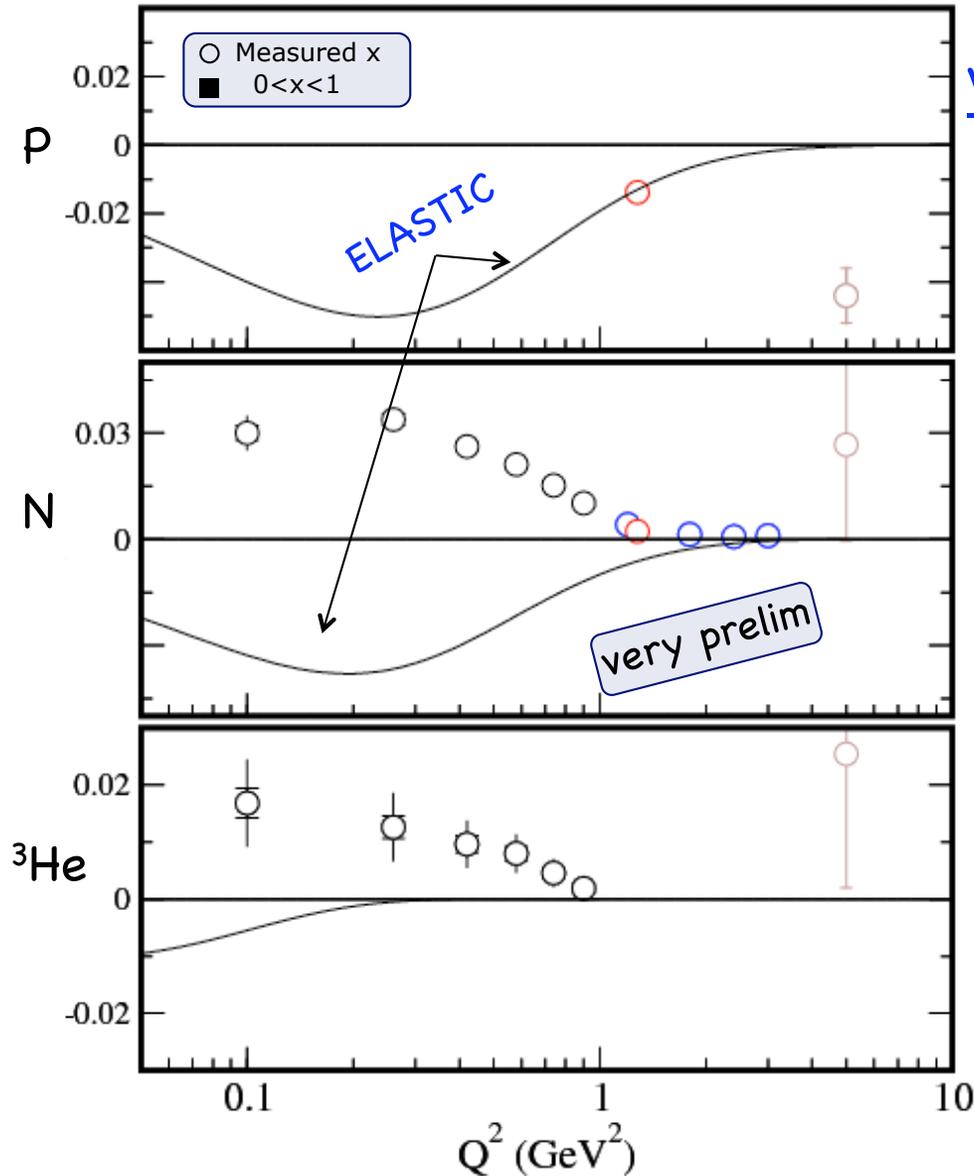
**Sane:** setting up now!

$$2.3 < Q^2 < 6 \text{ GeV}^2$$

**"g<sub>2p</sub>"** in Hall A, 2011

$$0.015 < Q^2 < 0.4 \text{ GeV}^2$$

# BC Sum Rule



## What can BC tell us about Low-X?

Alternatively, if we assume BC holds we can learn something about the unmeasured part of Integral

$$\int_0^1 g_2(x, Q^2) dx = 0$$

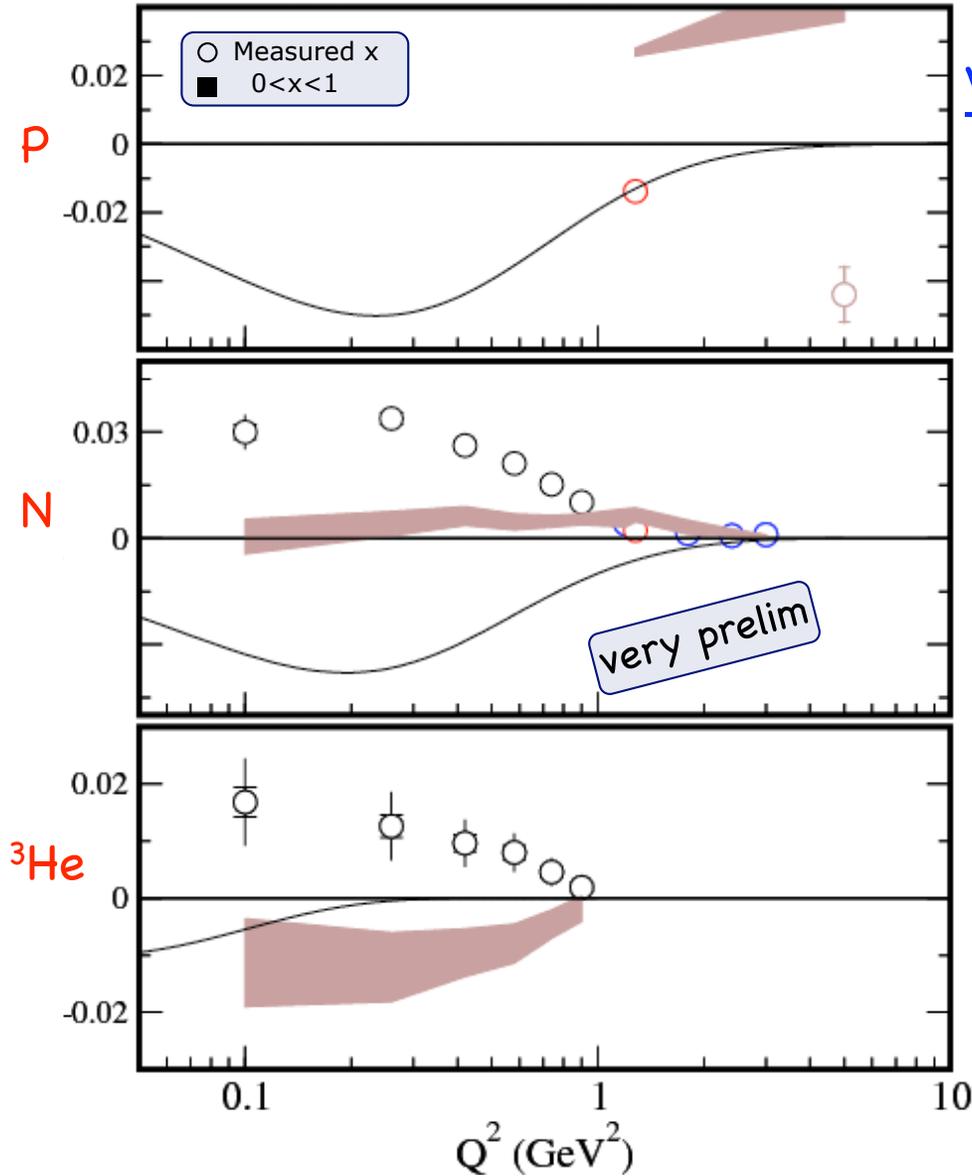
BC = 0 => Res + Elas + "DIS"

Unmeasured Low-x part

"DIS" = -(RES+ELAS)

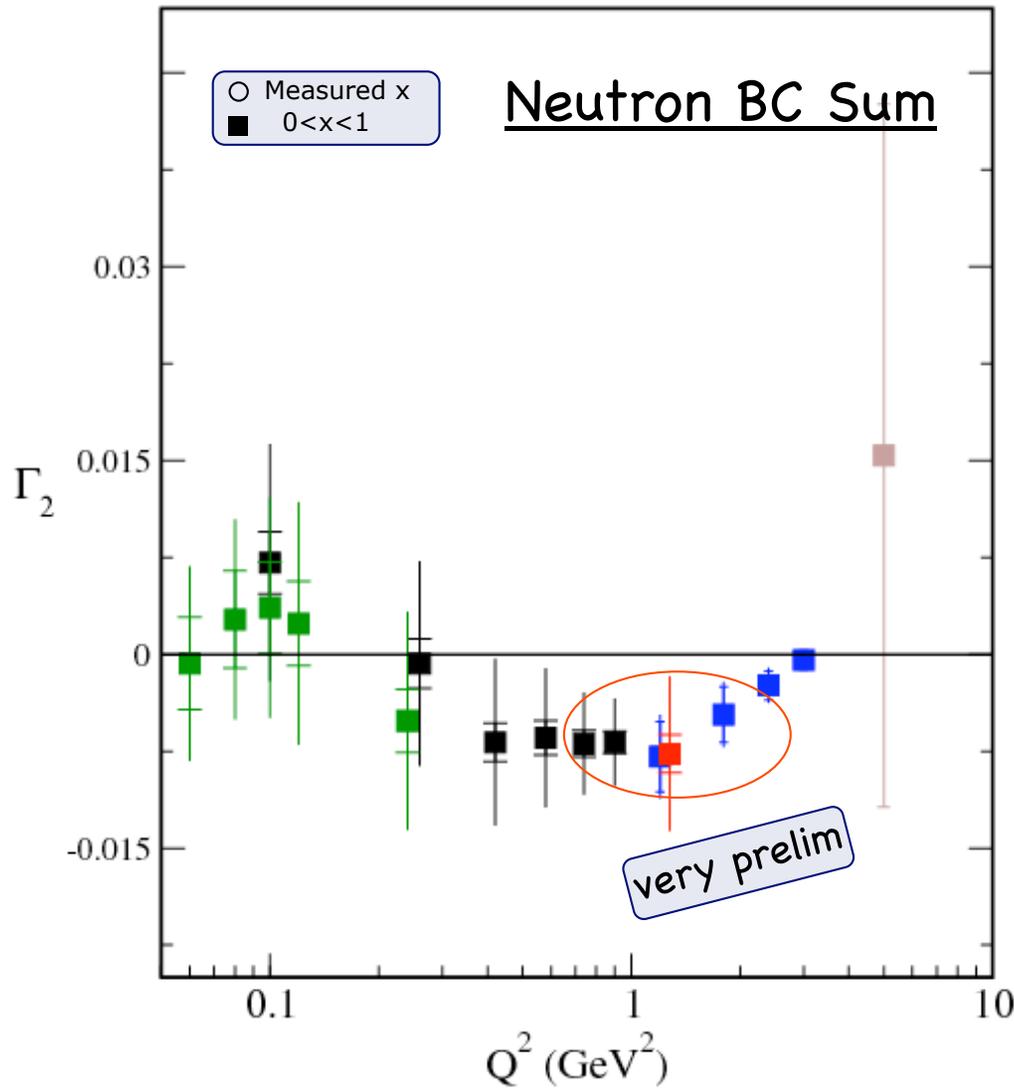
# BC Sum Rule

What can BC tell us about Low-X?



Unmeasured Low-X  
= -(RES+ELAS)

# BC Sum Rule



## Standard Deviations from Zero

$Q^2$	Tot $\sigma$	Stat $\sigma$
0.74	1.7	6.6
0.9	2.0	7.9
1.2	2.4	2.9
1.8	1.9	2.2
2.4	1.8	2.2

Just on the edge of being interesting

Statistical precision allows unambiguous test, but limited by large systematics.

Highly desirable to revisit E94010 systematics. RC and DIS could perhaps be improved with newer data.

# Higher Moment

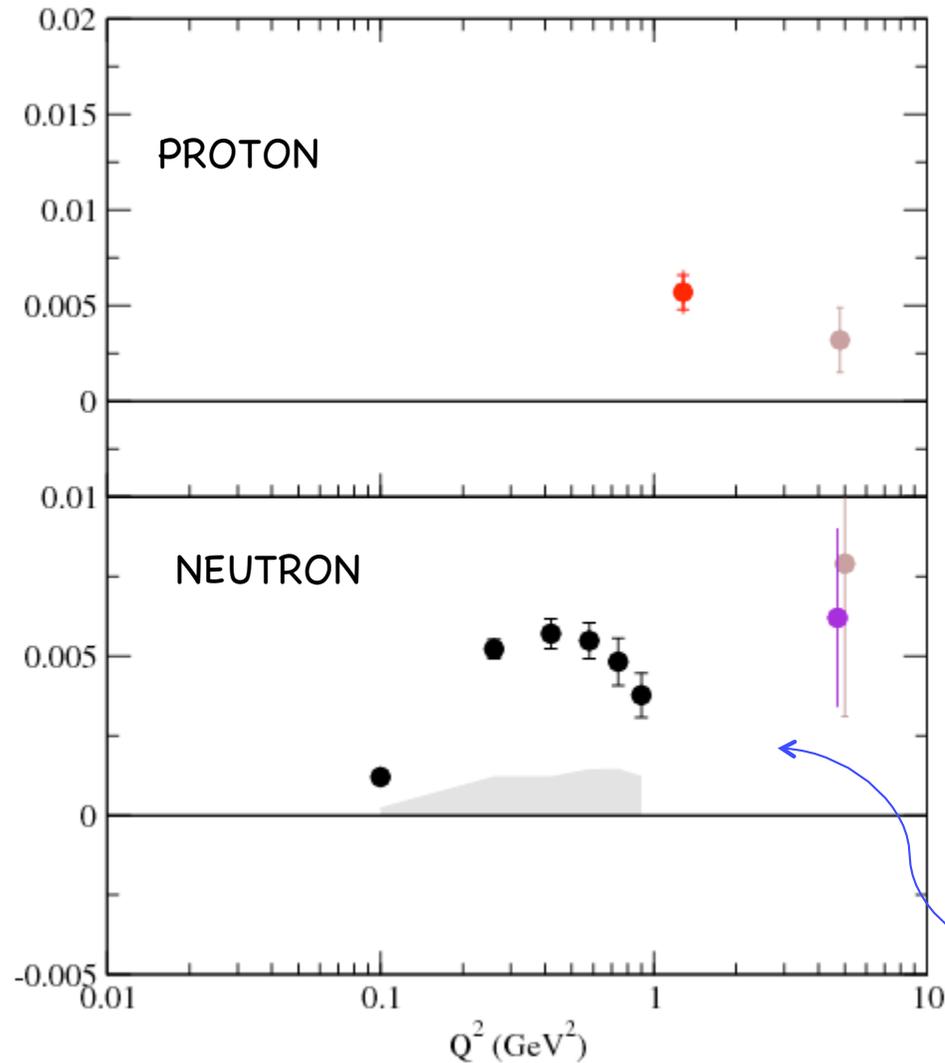
Cornwall Norton Moment

$$I(Q^2) = 2 \int_0^{1-\varepsilon} x^2 (2g_1 + 3g_2) dx$$

$I(Q^2) \neq$  the twist-3 matrix element but very interesting all the same.

More on this later...

# $I(Q^2)$



## Existing World Data on $I(Q^2)$

**BLACK** : E94010

M. Amarian, *et al.* PRL. 92 (2004) 022301

**BROWN** : E155

P.Anthony, *et al.* PLB. 553 (2003) 18

**RED** : RSS.

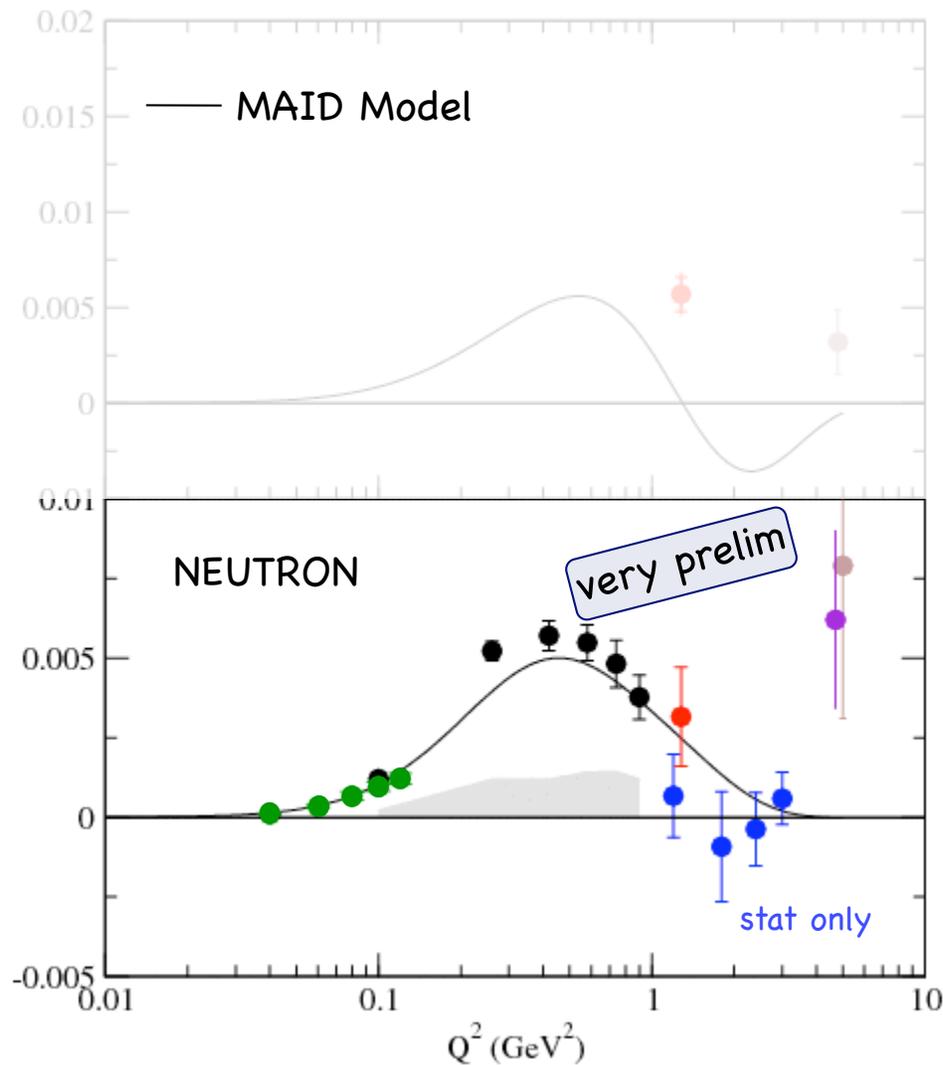
Wesselman, Slifer, Tajima *et al.* PRL 98(2007)132003.

**Magenta**: E99-117

X. Zheng *et al.* PRC 70(2004)065207

What's happening  
at large  $Q^2$ ?

# $I(Q^2)$



## BRAND NEW DATA!

### Very Preliminary

RED : RSS. (Hall C, NH<sub>3</sub>,ND<sub>3</sub>)

K. Slifer, O. Rondon *et al.* in preparation

BLUE: E01-012. (Hall A, <sup>3</sup>He)

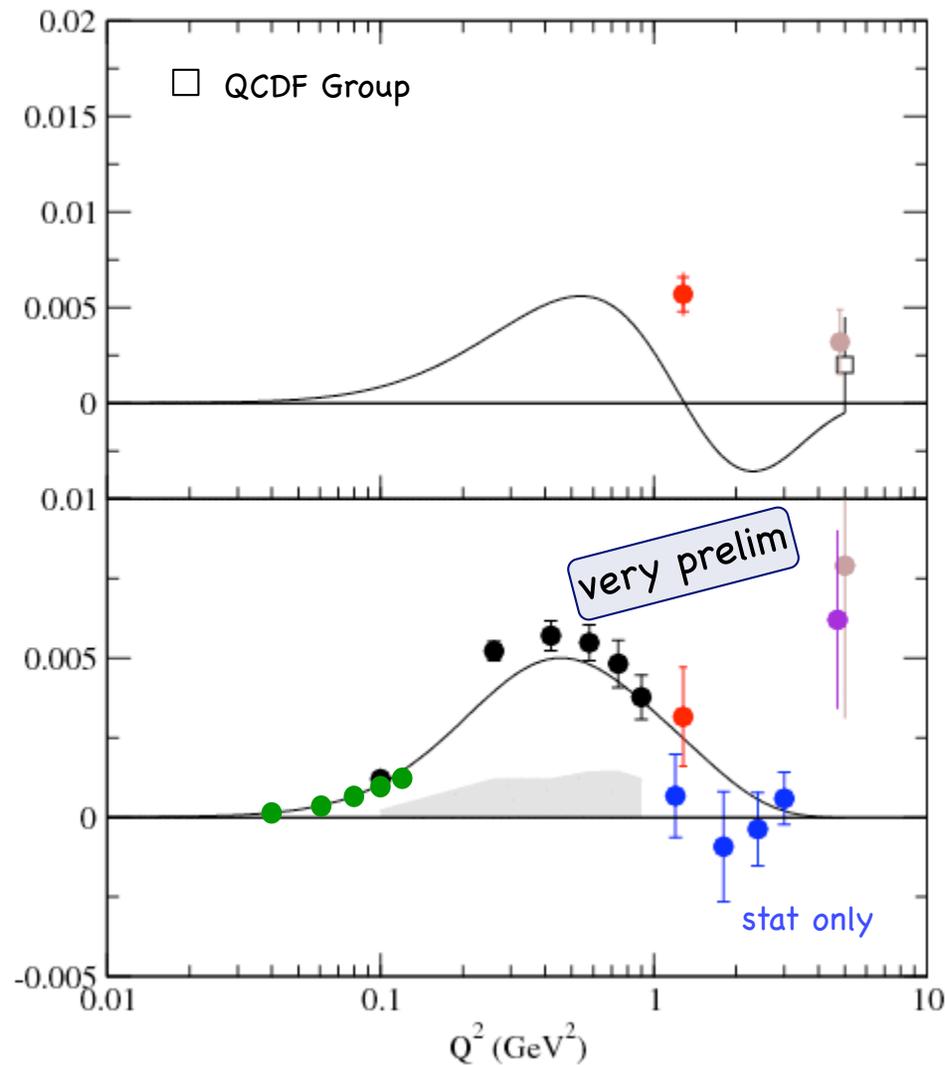
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GREEN: E97-110. (Hall A, <sup>3</sup>He)

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Thanks also to the spokesmen of these experiments!  
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# $I(Q^2)$



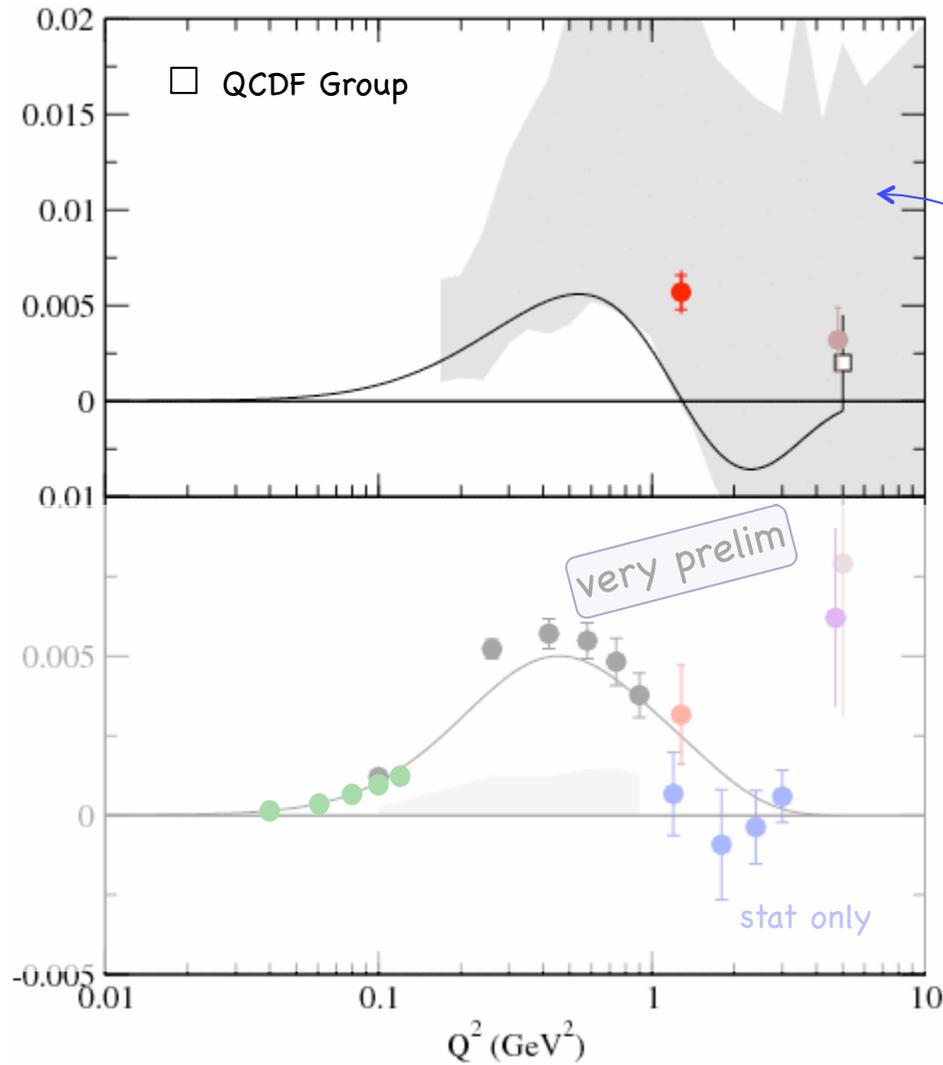
$I(Q^2) \rightarrow 0$  for  $Q^2=0$  and  $Q^2=\infty$

largest around  $Q^2=1$

Not too useful to the OPE in this region, but excellent gauge of "QCD complexity"

(i.e. where's the most difficult place to make any meaningful QCD calculation?)

# $I(Q^2)$

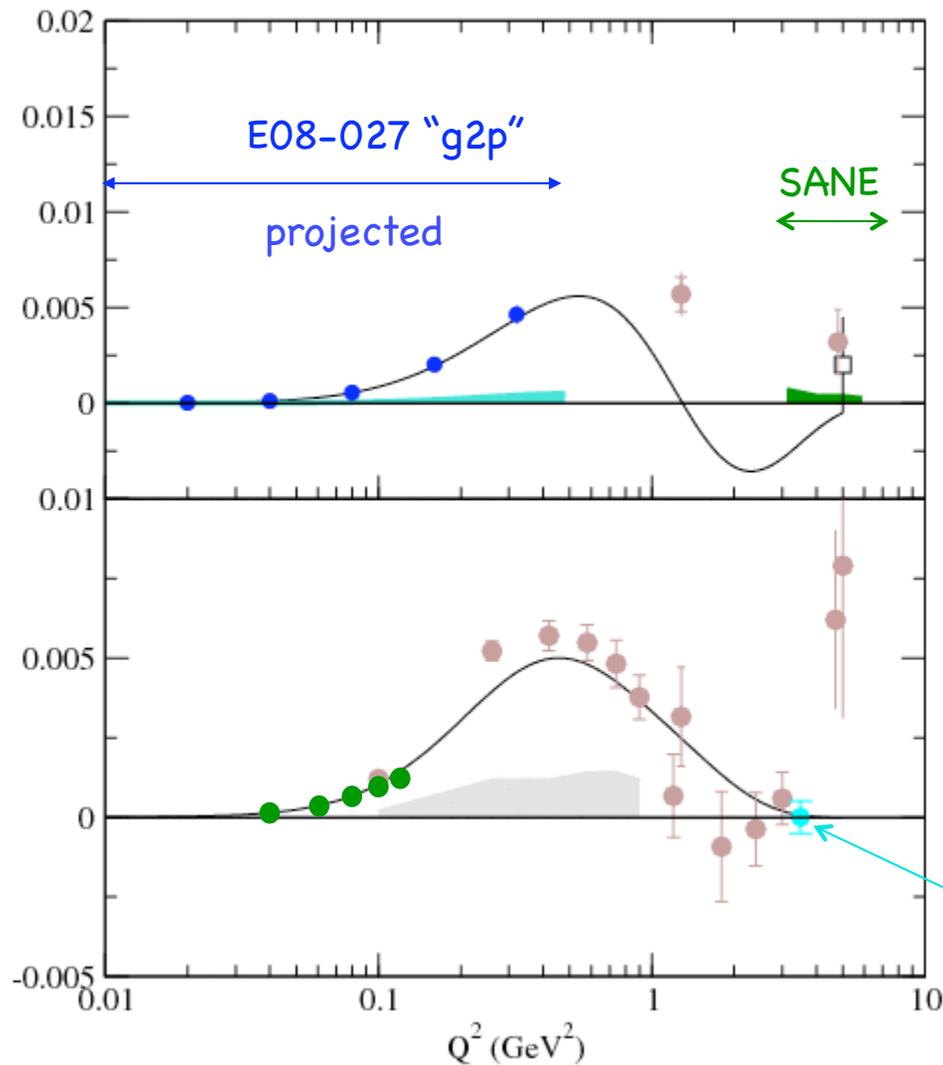


Osipenko *et al.* PRD. 71 (2005) 054007

Shaded Region: Estimate of  $I(Q^2)$

Large uncertainty due to lack of knowledge of  $g_2p$

# $I(Q^2)$



Upcoming 6 GeV Experiments

Sane: Fall 2008

$2.3 < Q^2 < 6 \text{ GeV}^2$

"g2p" in Hall A, 2011

$0.015 < Q^2 < 0.4 \text{ GeV}^2$

"d2n" in Hall A, 2009

$Q^2 = 3 \text{ GeV}^2$

# Operator Product Expansion

## Expansion of SF moments in powers of $1/Q^2$ ("twist")

example:

$$\Gamma_1(Q^2) = \int_0^1 g_1(x, Q^2) dx = \sum_{\tau=2,4,\dots} \frac{\mu_\tau(Q^2)}{Q^{\tau-2}}$$

$$\mu_4 = \frac{1}{9} M^2 (\tilde{a}_2 + 4\tilde{d}_2 + 4\tilde{f}_2)$$

leading twist      twist-3      twist-4

Lowest order (twist-2) maps to the successful parts of the parton model.

Higher twists arise from non-perturbative multiparton interactions

# Cornwall-Norton Moments

$$I(Q^2) = 2 \int_0^{1-\varepsilon} x^2 (2g_1 + 3g_2) dx$$
$$= \tilde{d}_2(Q^2) + \vartheta \left( \frac{M^2}{Q^2} \right)$$

Typical method of extracting  
twist-3 matrix element

But completely ignores TMC!

Very significant below  $Q^2 \approx 5$

Y.B. Dong PRC 77(2008) 015201

Y.B.Dong PLB 653,(2007)18

# Nachtmann Moments

Nachtmann Moments:

$$M_2^3(Q^2) = \int_0^1 dx \left( \frac{\xi^4}{x^2} \right) \left\{ \frac{x}{\xi} g_1 + \left[ \frac{3}{2} \left( \frac{x}{\xi} \right)^2 - \frac{3}{4} \frac{M^2}{Q^2} x^2 \right] g_2 \right\}$$
$$= \frac{\tilde{d}_2}{2}$$

Matsuda & Uematsu, N.Phys. B53(1998)301  
Piccione & Ridolfi N. Phys. B513(1998)301

Generalization of CN moments to protect from the TMC

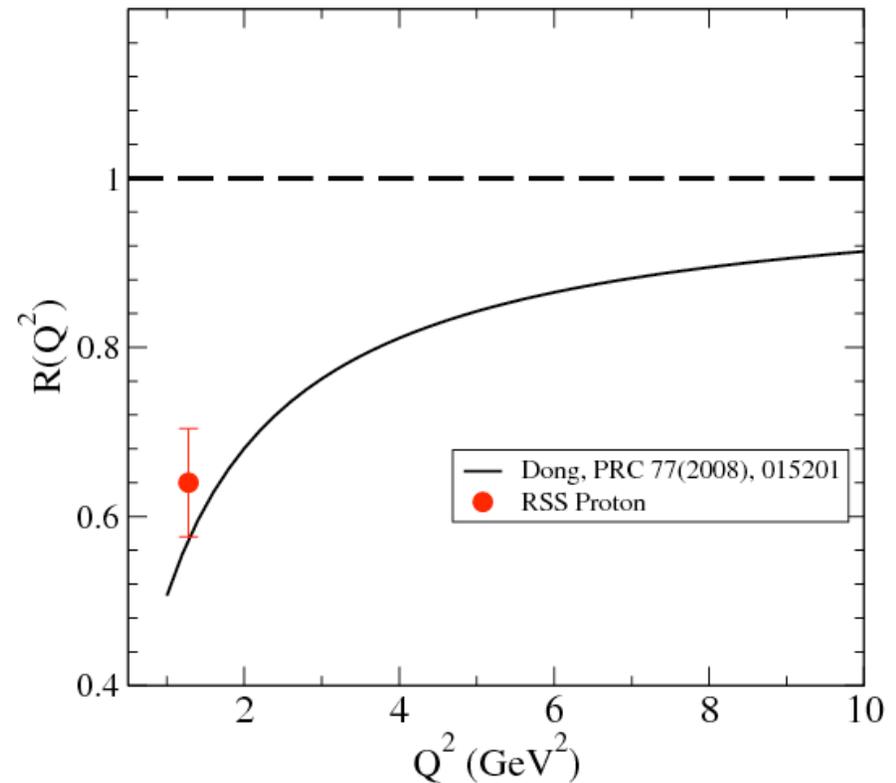
$$\frac{M^2}{Q^2} \rightarrow 0 \quad M_2^3 \rightarrow \int x^2 (2g_1 + 3g_2) dx \quad \text{Reduces to familiar form}$$

Not a new idea, but difficult to implement unless  $g_2$  measured simultaneously with  $g_1$

# Quantifying Size of TMC

$$R(Q^2) = \frac{2M_2^3(Q^2)}{I(Q^2)}$$

$R \rightarrow 1$  in case of vanishing nucleon mass



$R$  always less than 1  $\Rightarrow I(Q^2)$  overestimates twist-3

Target Mass Corrections must be applied in order to obtain clean dynamical Twist-3

# Generalization of $\Gamma_1$

$$M_1^1(Q^2) = \int_0^1 dx \left( \frac{x}{\xi} \right)^2 \left\{ \left[ \frac{x}{\xi} - \frac{1}{9} \left( \frac{M}{Q} \right)^2 x \xi \right] g_1 - \left( \frac{M}{Q} \right)^2 x^2 \frac{4}{3} g_2 \right\}$$

$$\frac{M^2}{Q^2} \rightarrow 0 \quad M_1^1 \rightarrow \Gamma_1$$

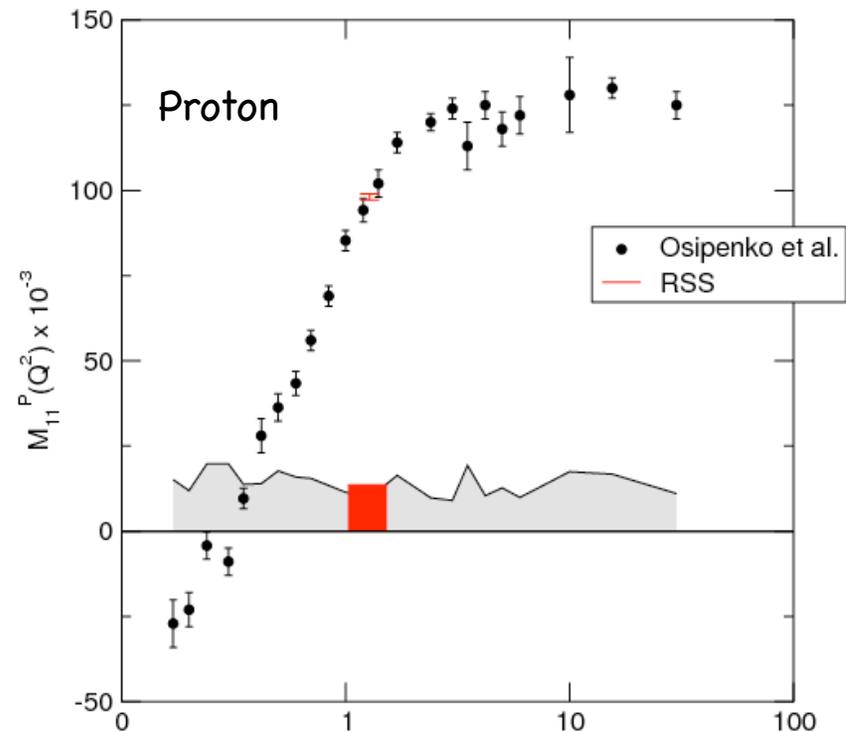
Matsuda & Uematsu, N.Phys. B53(1998)301

Piccione & Ridolfi N. Phys. B513(1998)301

Osipenko *et al.* PRD 71, 054007 (2005)

Global analysis of g1p data. Allows to cleanly extract leading twist term.

TMC not as large as for I(Q<sup>2</sup>)



# Summary

## Burkhardt–Cottingham Sum Rule

Good coverage for Neutron. Proton  $g_{2p}$  is still relatively unknown.

Data seems to validate BC, but at the  $2.5\sigma$  level around  $Q^2=1$   
Important to update the systematics of the old experiments

Assuming BC holds, we can use JLab data to say something about low- $x$ .

## Target Mass Effects

TMC are significant at JLab kinematics

Nachtmann moments protect the SSF from TMC

Must use Nachtmann Moments in order to cleanly extract Higher twists

## JLab 6 GeV Program

Still lots of Good Physics to be completed before the upgrade.

Backups

# References

## BLACK : E94010. (Hall A, $^3\text{He}$ )

M. Amarian, *et al.* PRL. 92 (2004) 022301  
K. Slifer, *et al.* PRL. 101:022303,2008

## RED : RSS. (Hall C, $\text{NH}_3, \text{ND}_3$ )

Wesselman, Slifer, Tajima *et al.*  
PRL 98(2007)132003.  
Slifer, Rondon *et al.* in preparation

## BROWN : E155. (SLAC $\text{NH}_3, ^6\text{LiD}$ )

P.Anthony, *et al.* PLB. 553 (2003) 18

## Magenta E99-117(Hall A, $^3\text{He}$ )

X. Zheng *et al.* PRC 70(2004)065207

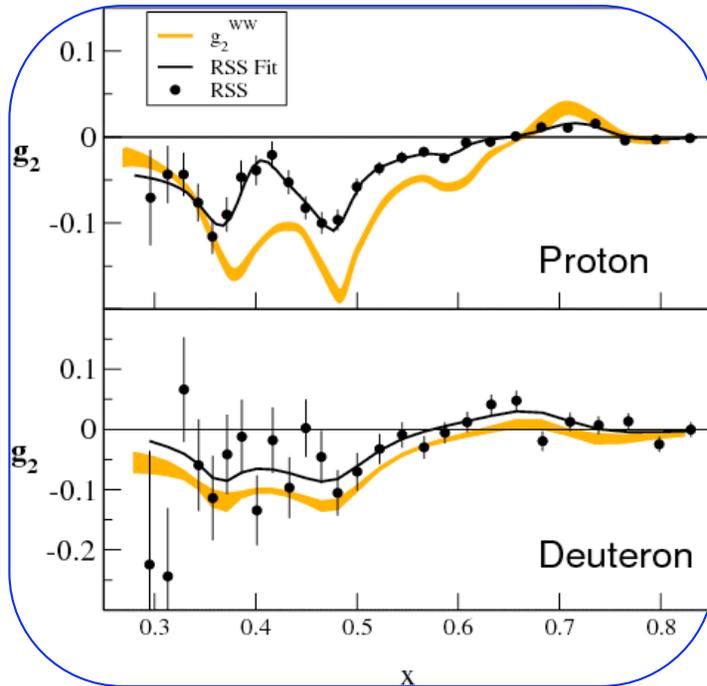
## BLUE: E01-012. (Hall A, $^3\text{He}$ )

P. Solvignon *et al.* arXiv:0803.3845 (PRL accepted)  
P. Solvignon *et al.* in preparation

## SHADED : Theory

Osipenko *et al.* PRD. 71 (2005) 054007

# BC Sum Rule



## Unmeasured Contributions

### Low-X Estimate

Assume  $g_2 = g_2^{ww}$  at low  $x$ .

Supported by RSS data

15% variation seen depending on choice of  $g_1$  used.

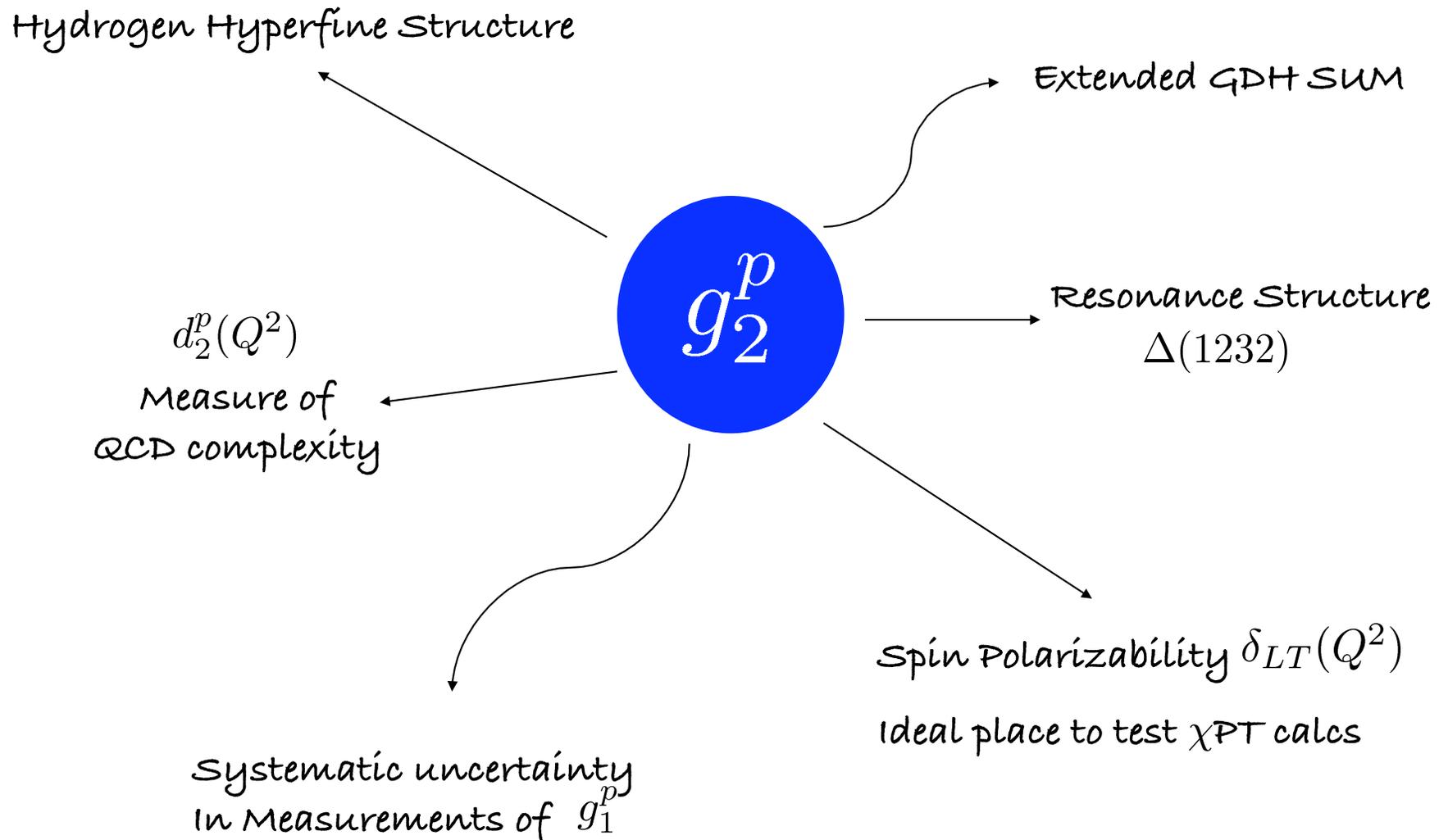
### ELASTIC: X=1

$$g_1^{el}(x, Q^2) = \delta(x-1) G_M(Q^2) \frac{G_E(Q^2) + \tau G_M(Q^2)}{2(1+\tau)}$$

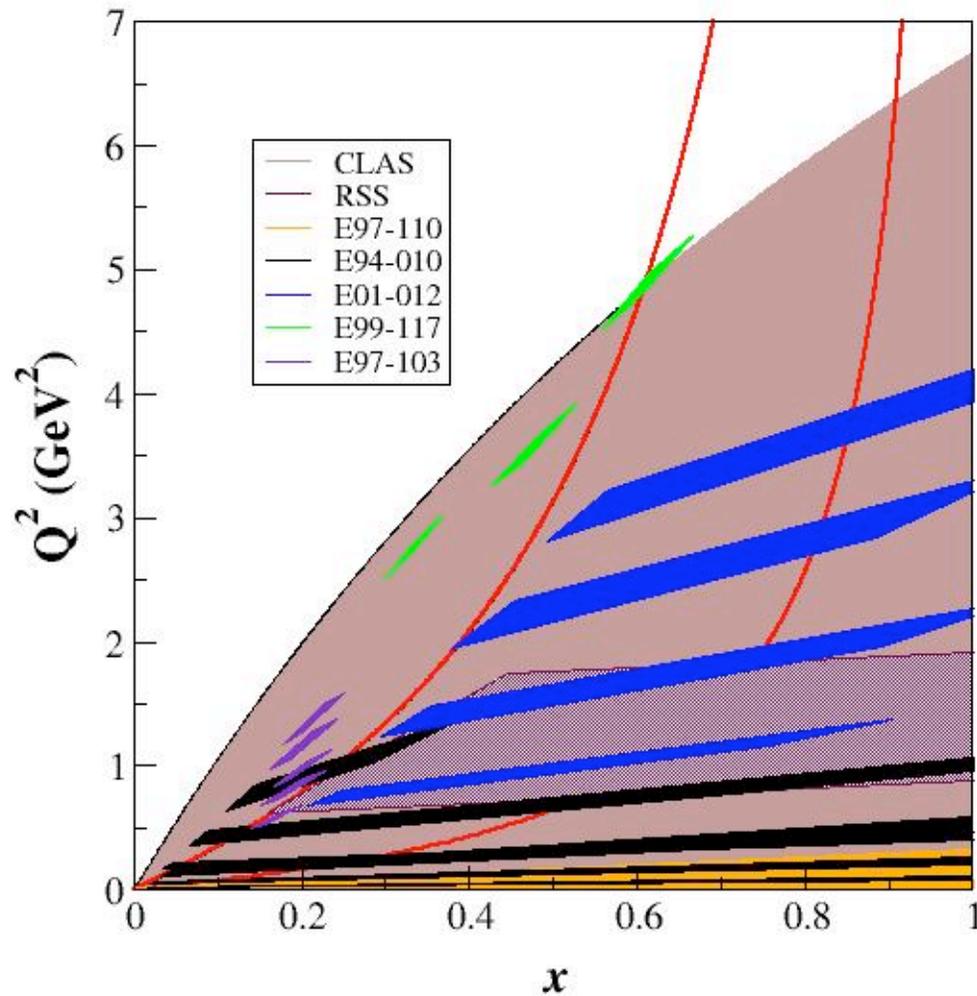
$$g_2^{el}(x, Q^2) = \delta(x-1) \tau G_M(Q^2) \frac{G_E(Q^2) - G_M(Q^2)}{2(1+\tau)}$$

(Form Factor uncertainties less than 5%)

# Summary



# JLab Kinematic Coverage



Overview of available kinematic range at JLab

Uniquely positioned to provide data in transition region of QCD

# Target Mass Corrections

Purely kinematic effects from finite value of  $4M^2x^2/Q^2$

$$g_1(x, Q^2) = g_1(x, Q^2, M = 0)$$

From PQCD

$$+ \frac{M}{Q^2} g_1^{(1)TMC}(x, Q^2)$$

Purely kinematical

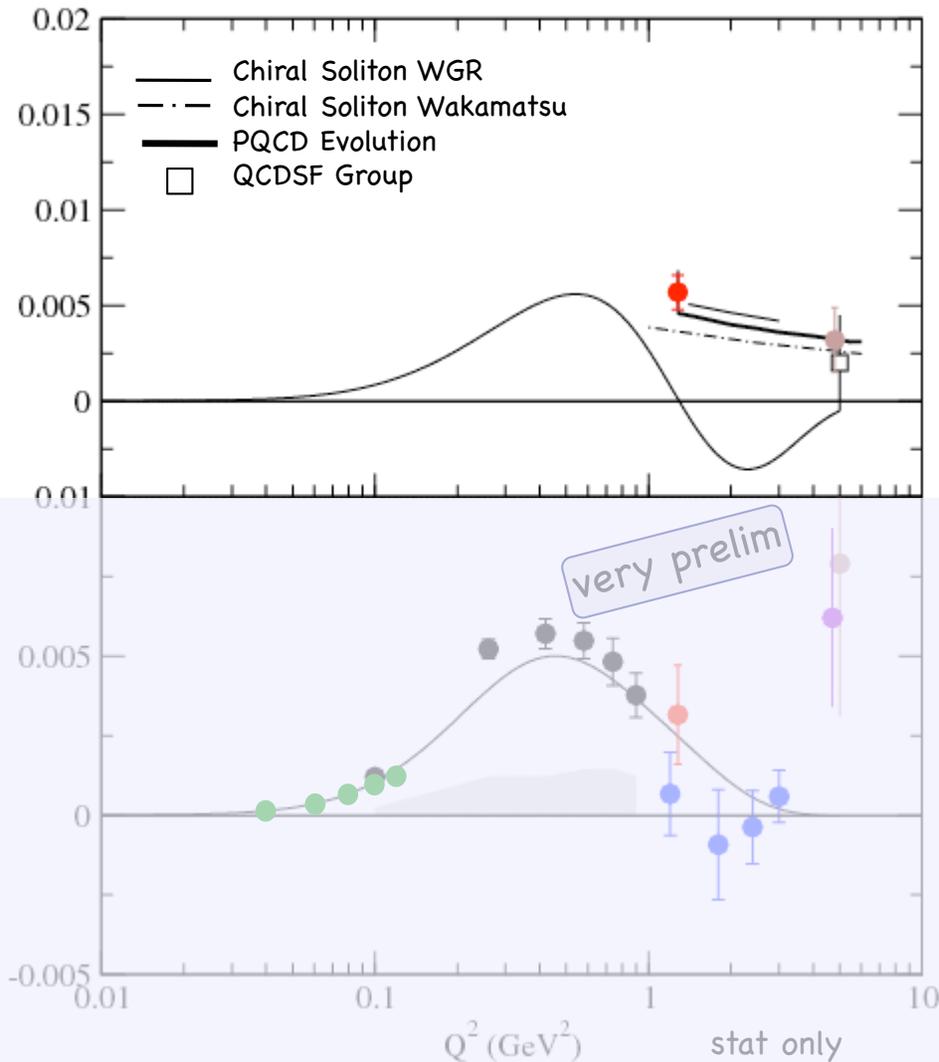
$$+ \frac{h(x, Q^2)}{Q^2} + \mathcal{O}(1/Q^4)$$

Higher twist

$$\int_0^1 x^2 g_1(x, Q^2) dx = \frac{1}{2} \tilde{a}_2 + \mathcal{O}\left(\frac{M^2}{Q^2}\right)$$

$$\int_0^1 x^2 g_2(x, Q^2) dx = \frac{1}{3} (\tilde{d}_2 - \tilde{a}_2) + \mathcal{O}\left(\frac{M^2}{Q^2}\right)$$

# I(Q<sup>2</sup>)



$Q^2$  evolution predicted well by PQCD and Chiral Soliton models.

WGR: PRD55 (1997) 6910

Wakamatsu: PLB 487(2000)118

PQCD: Nucl. Phys. B201 (1982) 141

QCDSF: PRD63, 074506(2001)

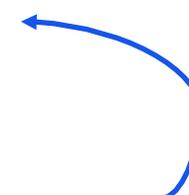
# $g_2$ Structure Function

Wandzura-Wilczek relation

PLB 72 (1977) 195

$$g_2^{WW}(x, Q^2) = -g_1(x, Q^2) + \int_x^1 \frac{g_1(y, Q^2)}{y} dy$$

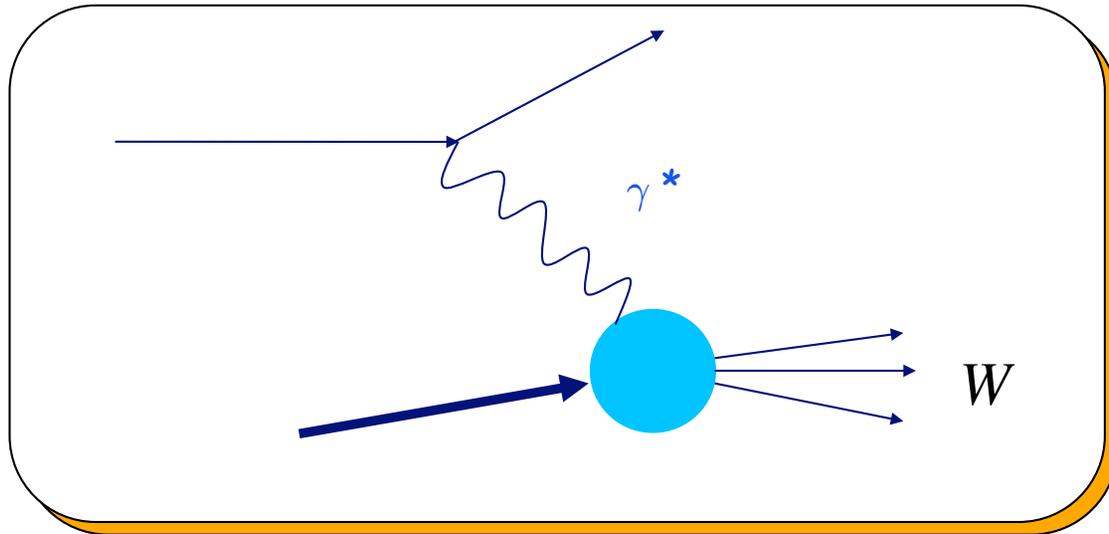
Leading twist determined entirely by  $g_1$

$$g_2 = g_2^{WW} + \bar{g}_2$$


Higher twist

$g_2$  doesn't exist in Parton Model.  
Good quantity to study higher twist

# Inclusive Electron Scattering



$$\frac{d^2\sigma}{d\Omega dE'} = \sigma_{Mott} \left[ \frac{1}{\nu} F_2(x, Q^2) + \frac{2}{M} F_1(x, Q^2) \tan^2 \frac{\theta}{2} \right]$$

Inclusive Cross Section

deviation from point-like behavior  
characterized by the **Structure Functions**