



"Beam in 30 minutes or it's free"

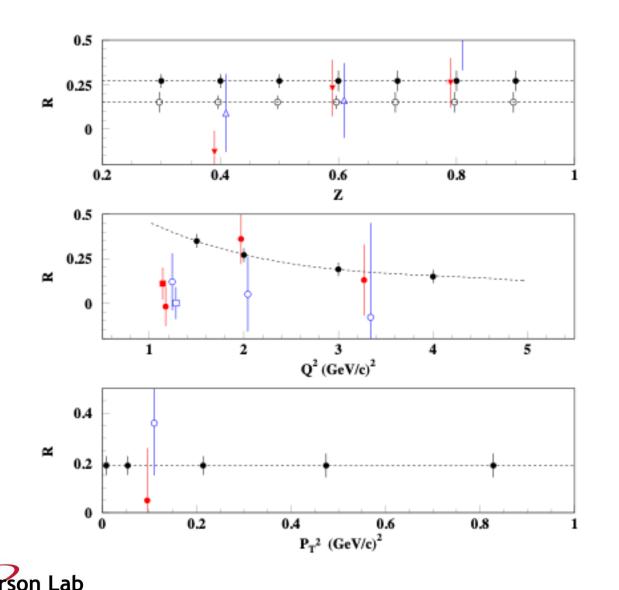
#### February 3, 2025

Topics:

- 1. Analysis goals
- 2. Rosenbluth Separations
- 3. Extraction of cross sections
- 4. Other stuff



#### **R-SIDIS Measurement Goals**



Extract  $R = \sigma_L / \sigma_T$  in SIDIS  $\rightarrow$  charged pions and kaons  $\rightarrow$  H and D targets

Measure as a function of x/Q<sup>2</sup>, z, and  $\mathsf{P}_\mathsf{T}$ 

Compare R in SIDIS to DIS Key questions:

- How does R transition from low z to exclusive limit (z=1)?
- 2. Is R the same for pions and kaons?
- 3. What is the  $P_T$  dependence?

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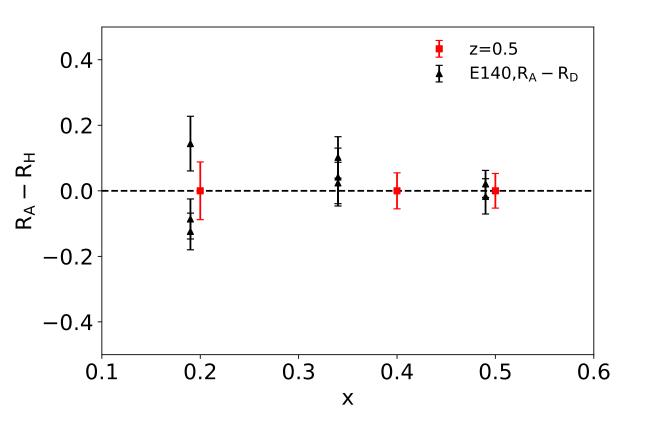
## **R-SIDIS Measurement Goals**

Extract nuclear dependence of  $R_{SIDIS}$ 

- → Approved experiment to measure nuclear dependence of R in inclusive DIS – related to EMC effect, anti-shadowing, etc.
- → Nuclear dependence of R<sub>SIDIS</sub> important for interpretation of hadron attenuation measurements, dilution extraction for polarized target experiments, possible interesting new physics

Exploratory measurement  $\rightarrow x/Q^2$  dependence, z dependence, P<sub>T</sub> dependence at subset of proton/deuteron measurements

SLAC E140: Nuclear Dependence of R in DIS
 E12-24-001: Nuclear Dependence of R in SIDIS (projected precision)





#### Inclusive Electron Scattering cross section and kinematics

 $\frac{d\sigma}{d\Omega dE'} = \frac{4\alpha^2 (E')^2}{Q^4} \left[ W_2(\nu, Q^2) \cos^2 \frac{\theta}{2} + 2W_1(\nu, Q^2) \sin^2 \frac{\theta}{2} \right] \qquad \frac{MW_1(n, Q^2) \ \mathbb{R} \ F_1(x)}{nW_2(n, Q^2) \ \mathbb{R} \ F_2(x)} \qquad F_2(x) = \sum_i e_i^2 x q_i(x)$ 

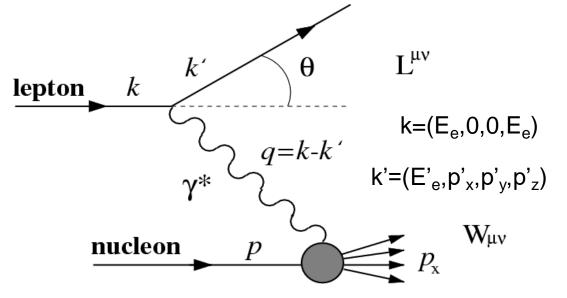
 $Q^2 = -(\text{four-momentum transferred to struck nucleon})^2 = 4E_e E'_e \sin^2 \frac{\theta_e}{\Omega}$ 

 $\nu = E_e - E'_e$ 

W<sup>2</sup>=total energy of virtual-photon + target in CM frame

$$= \nu^2 + M^2 + 2M\nu \ \ \text{-}\ \text{q}^2$$

 $x = \frac{Q^2}{2Mu}$  Bjorken scaling variable  $\Rightarrow$  Fraction of nucleon momentum carried by struck quark



## (Inclusive) Rosenbluth (L-T) Separations

Inclusive cross section can be re-written:

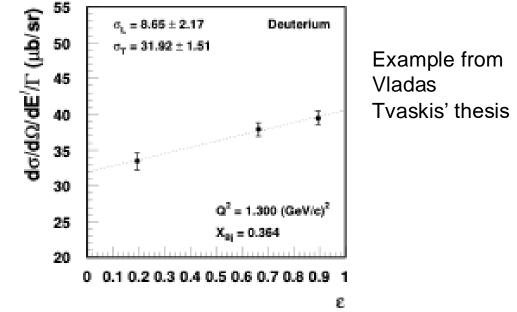
$$\frac{d\sigma}{d\Omega dE} = \Gamma \left[ \sigma_T(\nu, Q^2) + \epsilon \sigma_L(\nu, Q^2) \right]$$

$$\Gamma = \frac{\alpha}{2\pi^2} \frac{E'_e}{E_e} \frac{1}{Q^2} \frac{1}{1-\epsilon} K_{eq} \qquad K_{eq} = \frac{W^2 - M^2}{2M} \qquad \epsilon = \left[1 + 2\left(1 + \frac{Q^2}{4M^2x^2}\right) \tan^2\frac{\theta}{2}\right]^{-1}$$

Plot cross section at fixed v (or x) and Q<sup>2</sup>  $\rightarrow$  Plot vs.  $\varepsilon$  and fit line

 $\rightarrow$  Slope =  $\sigma_L$ , intercept  $\sigma_T$ 

Vary  $\epsilon$  by making measurements at different beam energies (scattered electron momentum and angle will also change to keep x and Q<sup>2</sup> fixed)

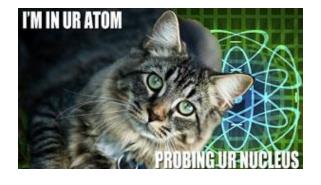




## Nuclear Dependence of $R = \sigma_L / \sigma_T$

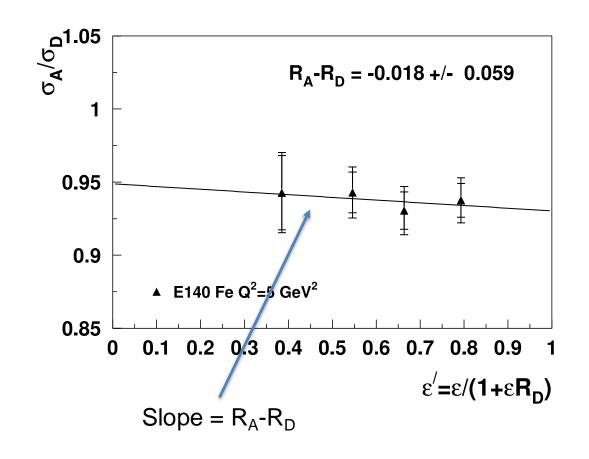
Nuclear dependence of R can be extracted by looking epsilon dependence of target ratios

$$\frac{\sigma_A}{\sigma_H} = \frac{\sigma_A^T}{\sigma_H^T} [1 + \epsilon' (R_A - R_H)]$$
$$\epsilon' = \epsilon / (1 + \epsilon R_H)$$



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Example: R<sub>A</sub>-R<sub>D</sub> from SLAC E140 (inclusive DIS)



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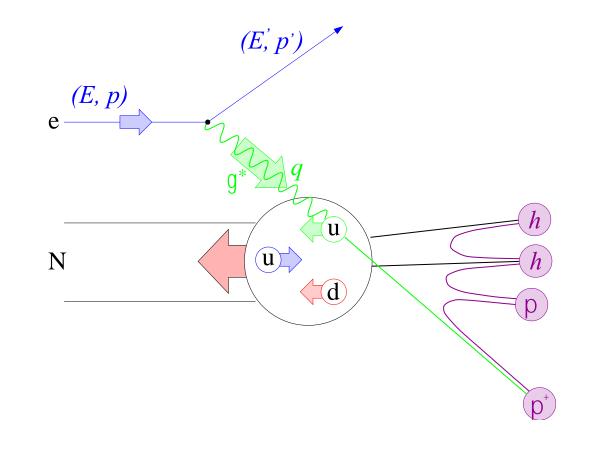
## **Semi-inclusive DIS**

SIDIS  $\rightarrow$  production of one or more hadrons in DIS reaction

Simple picture:

1. Electron scatters from quark in nucleon 2. Quark is kicked out  $\rightarrow$  subsequently hadronizes, ending up in bound state

In this simple picture, SIDIS can be used to "tag" the flavor of the struck quark in DIS process



$$d\sigma^h \propto \sum f^{H 
ightarrow q}(x) \mathrm{d}\sigma_q(y) D^{q 
ightarrow h}(z)$$

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Fragmentation function

#### **Semi-inclusive DIS**

In principle SIDIS can also be used to gain information about spatial distributions of quar in nucleons

→ This requires measurements/observation c transverse degrees of freedom

DIS can also be used to gain  
bout spatial distributions of quarks  
es measurements/observation of  
agrees of freedom  
$$d\sigma^{h} \propto \sum f^{H \to q}(x) d\sigma_{q}(y) D^{q \to h}(z)$$
$$\downarrow$$
$$d\sigma^{h} \propto \sum f^{H \to q}(x, k_{T}) \otimes d\sigma_{q}(y) \otimes D^{q \to h}(z, p_{\perp})$$



## **SIDIS Cross Section(s)**

Unpolarized cross section  $\rightarrow$  requires 3 more degrees of freedom

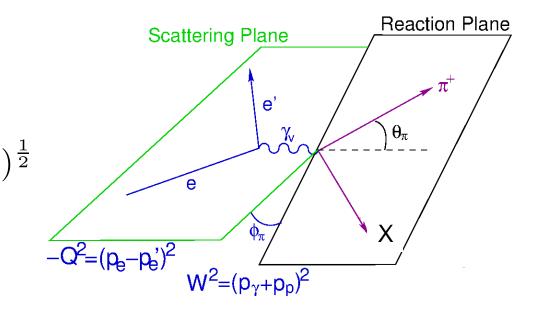
$$\frac{d\sigma}{dE_e d\Omega_e dP_\pi d\Omega_\pi} = \Gamma \left[ \frac{d\sigma_T}{dP_\pi d\Omega_\pi} + \epsilon \frac{d\sigma_L}{dP_\pi} d\Omega_\pi + \sqrt{2\epsilon(1+\epsilon)} \frac{d\sigma_{LT}}{dP_\pi d\Omega_\pi} \cos\phi + \epsilon \frac{d\sigma_{TT}}{dP_\pi d\Omega_\pi} \cos 2\phi \right]$$
$$\frac{d\sigma}{dx dy dz dp_T^2 d\phi} = \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\epsilon)} \left[ F_T + \epsilon F_L + \sqrt{2\epsilon(1+\epsilon)} \cos\phi F_{LT} + \epsilon \cos 2\phi F_{TT} \right]$$

Inclusive: structure functions depend just on x and  $Q^2$ SIDIS: structure functions depend on x,Q<sup>2</sup>, z, and P<sub>T</sub>

$$z = \frac{q \cdot p}{q \cdot P} = \frac{E_h}{\nu} \qquad p_{\parallel} = \frac{p \cdot q}{|q|} \qquad p_T = (p^2 - p_{\parallel}^2)$$

$$\cos \phi = \frac{(-\vec{q} \times \vec{k}) \cdot (-\vec{q} \times \vec{p_h})}{|\vec{q} \times \vec{k}| |\vec{q} \times \vec{p_h}|}$$

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#### **Rosenbluth separations for SIDIS**

$$\frac{d\sigma}{dE_e d\Omega_e dP_\pi d\Omega_\pi} = \Gamma \left[ \frac{d\sigma_T}{dP_\pi d\Omega_\pi} + \epsilon \frac{d\sigma_L}{dP_\pi} d\Omega_\pi + \sqrt{2\epsilon(1+\epsilon)} \frac{d\sigma_{LT}}{dP_\pi d\Omega_\pi} \cos\phi + \epsilon \frac{d\sigma_{TT}}{dP_\pi d\Omega_\pi} \cos 2\phi \right]$$

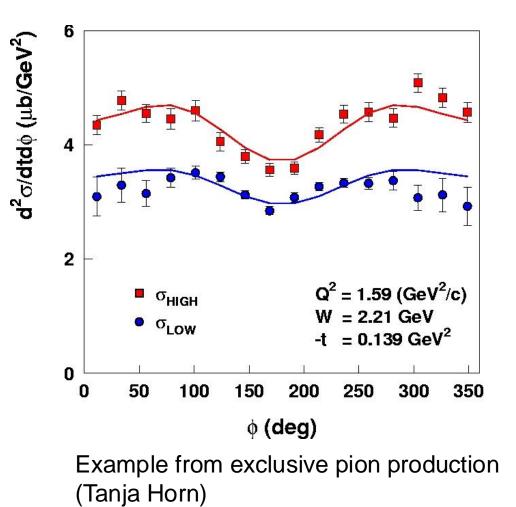
L-T separation for meson production more complicated due to phi dependent terms

We're most interested in  $\sigma_{\text{L}}$  and  $\sigma_{\text{T}}$ 

Need to extract cross section at fixed x, Q<sup>2</sup>, z, and P<sub>T</sub> vs.  $\phi$  at each beam energy ( $\epsilon$ )

Two options:

- 1. Integrate over  $\phi$  at each  $\varepsilon \rightarrow$  cross section reduces to  $\sigma_L + \varepsilon \sigma_T$
- 2. Do multiparameter fit over all  $\varepsilon$  settings  $\rightarrow$  extract L, T, LT, TT terms simultaneously





## Phase Space at high and low epsilon

Complication in L-T separations due to finite acceptance of spectrometers

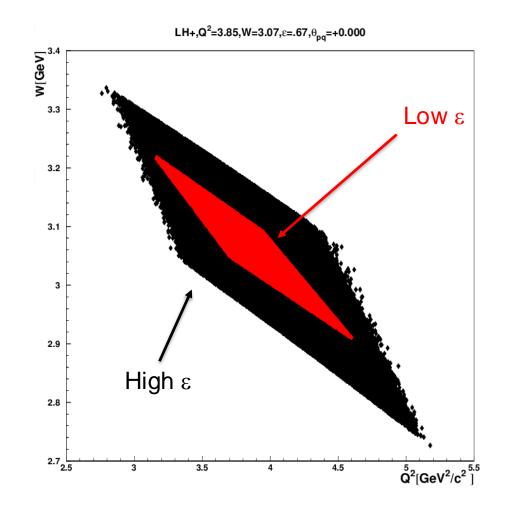
→ Electron phase space (x,Q<sup>2</sup>) or (W,Q<sup>2</sup>)

Need to perform L-T separation at same x and Q<sup>2</sup>

→ Integrating over different regions can results in different effective x and Q<sup>2</sup>

2 options:

- Add so-called "diamond cuts" to acceptance → force x/Q<sup>2</sup> at coverage low and high ε to be the same
  - Results in reduced statistics for high  $\epsilon$  data (or longer run times)
- "Bin-center" over full x/Q<sup>2</sup>
  - Bin-centering at some level is required anyway
  - Can result in larger systematic uncertainties





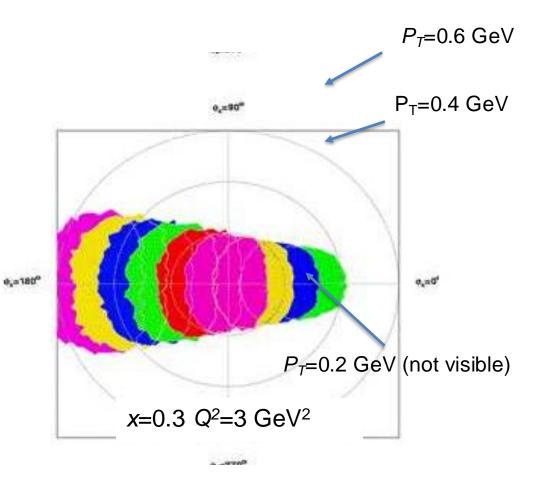
#### $P_T$ range and $\phi$ dependence

$$\frac{d\sigma}{dE_e d\Omega_e dP_\pi d\Omega_\pi} = \Gamma \left[ \frac{d\sigma_T}{dP_\pi d\Omega_\pi} + \epsilon \frac{d\sigma_L}{dP_\pi} d\Omega_\pi + \sqrt{2\epsilon(1+\epsilon)} \frac{d\sigma_{LT}}{dP_\pi d\Omega_\pi} \cos\phi + \epsilon \frac{d\sigma_{TT}}{dP_\pi d\Omega_\pi} \cos 2\phi \right]$$

The P<sub>T</sub> range over which we can unambiguously extract  $\sigma_L$  and  $\sigma_T$  is limited by the spectrometer acceptance

→Can move the spectrometer left and right of the q-vector direction to cover  $\phi=0$  and 180 degrees

→ Out-of-plane acceptance limits PT range near  $\phi$ =90 and 270 degrees





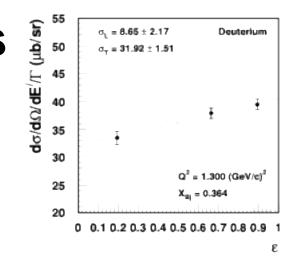
## **Systematic Uncertainties**

In general – there are 2 classes of systematic uncertainties

- Scale or normalization type systematic uncertainties
  - These uncertainties are in general independent of time and running condition
  - Example: target length → this will stay the same throughout run (unless damaged)
- Point-to-point or random systematic uncertainties
  - These uncertainties can vary with time, beam energy, spectrometer setting, etc.
  - Examples: acceptance (same region not populated at high and low epsilon) tracking efficiency (rate dependendent)

Point-to-point uncertainties are the most crucial to control

- Added in quadrature with statistical uncertainty at each epsilon
   → direct impact on slope, intercept
- Scale uncertainties will cancel when forming  $\sigma_{\text{L}}/\sigma_{\text{T}}$  ratio



	Type of systematic uncertainty		
	pt-to-pt	t-correlated	scale
Source	(%)	(%)	(%)
Acceptance	0.4	0.4	1.0
Target Thickness		0.2	0.8
Beam Charge		0.2	0.5
HMS+SHMS Tracking	0.1	0.1	1.5
Coincidence Blocking		0.2	
PID		0.4	
$\pi$ Decay	0.03		0.5
$\pi$ Absorption		0.1	1.5
Monte Carlo Generator	0.2	1.0	0.5
Radiative Corrections	0.1	0.4	2.0
Offsets	0.4	1.0	
Quadrature Sum	0.6	1.6	3.3
Fpi-2 Values	0.9	1.9	3.5



## **Corrections and backgrounds**

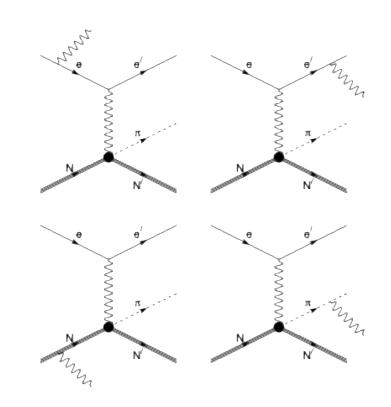
Radiative corrections  $\rightarrow$  emissions of "extra" photons by incoming/outgoing electrons, outgoing pion

Contributions:

- → SIDIS events from other kinematics can radiate into experimental acceptance → can also radiate out of acceptance
- → Events from exclusive processes can radiate into acceptance

→ For example: H(e,e' $\pi$ +)n, H(e,e' $\pi$ +) $\Delta^0$ 

Additional background possible from exclusive vector meson production  $\rightarrow$  can decay into  $\pi^+ \pi^-$  pair  $\rightarrow$  Dominate by diffractive  $\rho^0$  production





## **Corrections and backgrounds**

Exclusive

0.7

0.7

0.8

z

• π • π΄

0.8

z

Example from Hem Bhatt's thesis

0.2

0.1

n

0.2

0.1

0

son Lab

0.3

0.3

x=0.3

0.4

0.4

x=0.5  $Q^2 = 6.1 \text{ GeV}^2$ 

 $Q^2 = 3.5 \text{ GeV}^2$ 

0.5

0.5

0.6

0.6

γ(p→π<sup>+</sup>π<sup>-</sup>)γ(e′πX)

- $\rightarrow$  Exclusive, and delta production estimates based on models that describe the Hall C 2018 SIDIS and world exclusive data
- $\rightarrow$  Rho model based on HERMES modifications to Pythia, with tweaks to improve agreement with JLab 6 GeV results

0.8

z

0.8

z

Rho

0.7

0.7

۲(e/πX) 0.0 0.0 0.0

0.1

n

0.1

0.05

0.3

0.3

x=0.3

 $Q^2 = 3.5 \text{ GeV}^2$ 

0.4

 $Q^2 = 6.1 \text{ GeV}^2$ 

0.4

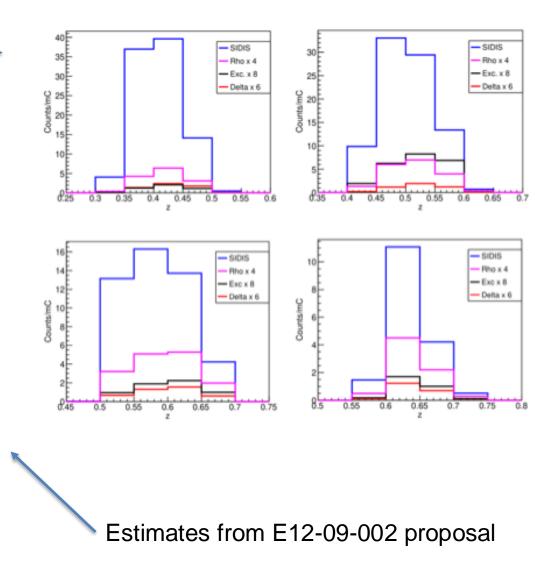
x=0.5

0.5

0.5

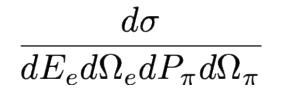
0.6

0.6



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#### **Extraction of cross sections**



# of scattered electrons in  $E_e, \Omega_e$  bin in coincidence with a pion in given  $P_{\pi}$ ,  $\Omega_{\pi}$  bin

 $\Delta E_e \Delta \Omega_e \Delta P_\pi \Delta \Omega_\pi$  (# of target protons/cm<sup>2</sup>) (# of incident electrons)

In principle, can calculate cross section "by-hand"

- $\rightarrow$  # of events from output of analyzer hcana
- $\rightarrow$  Incident electrons from beam current measurements
- $\rightarrow$  Target particle from target density and thickness
- $\rightarrow$  Phase space from Monte Carlo simulation

Also need efficiency corrections:

- $\rightarrow$  Live times (computer and electronic)
- → Tracking efficiency
- $\rightarrow$  Detector efficiencies

Other corrections:

- → Radiative effects
- $\rightarrow$  Bin centering
- $\rightarrow$  Pion decay



# **Bin centering**

When determining cross section, what kinematics do we quote?

→ Events are integrated over non-zero phasespace in (x,Q<sup>2</sup>)

Can we just average x and Q<sup>2</sup>?

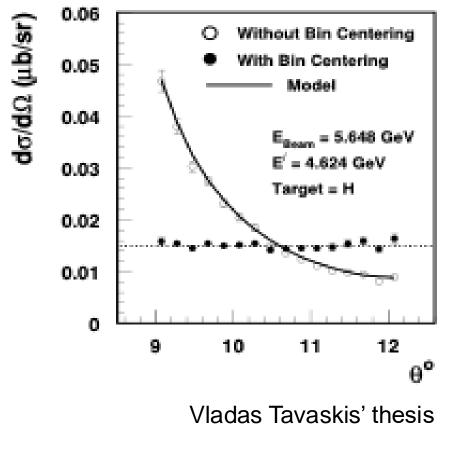
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- → If cross section is linear in those variables over the bin size, this could be ok
- → If cross section not linear average cross section at average (x,Q<sup>2</sup>) is not the same as the cross section evaluated at average (x,Q<sup>2</sup>)

Need to "bin-center" the data to quote cross-section at particular  $(x,Q^2)$ 

- → This can be done by applying explicit weight to each event OR
- → Can be done implicitly by comparing data to crosssection weighted simulation

#### Inclusive example



For coincidence reactions this is easier

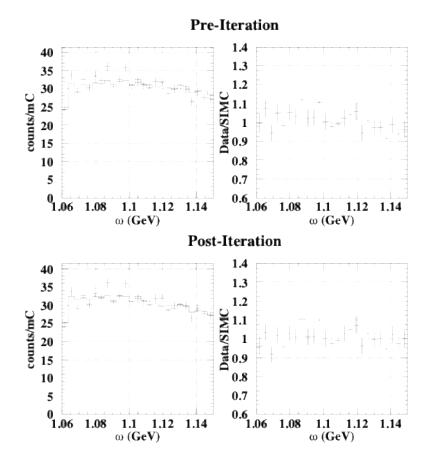
#### **Extraction of cross sections**

Easiest technique is the ratio method:

$$\left[\frac{d\sigma(x_0, Q_0^2, z_0, P_{T0})}{dE_e d\Omega_e dP_\pi d\Omega_\pi}\right]_{exp} = \frac{Y_{DATA}}{Y_{MC}(\sigma_{model})} \left[\frac{d\sigma(x_0, Q_0^2, z_0, P_{T0})}{dE_e d\Omega_e dP_\pi d\Omega_\pi}\right]_{model}$$

Monte Carlo includes radiative effects, pion decay, multiple scattering, acceptance ...

By using the ratio method is performing an "implicit" bin centering correction – but cross section model used in the MC must agree with the data  $\rightarrow$  may require correction and iteration of the model cross section



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#### **Extraction of cross sections - SIDIS**

$$\left[\frac{d\sigma(x_0, Q_0^2, z_0, P_{T0})}{dE_e d\Omega_e dP_\pi d\Omega_\pi}\right]_{exp} = \frac{Y_{DATA}}{Y_{MC}(\sigma_{model})} \left[\frac{d\sigma(x_0, Q_0^2, z_0, P_{T0})}{dE_e d\Omega_e dP_\pi d\Omega_\pi}\right]_{model}$$

Need to decide how you want to bin/present data, e.g., show z-dependence at fixed  $P_T$ ,  $P_T$  dependence at fixed z, etc.

 $Y_{data} \rightarrow$  efficiency-corrected # of counts normalized to charge

 $\rightarrow$  counts from root trees, charge and efficiencies from report file output from hcana

 $Y_{MC} \rightarrow$  Monte Carlo yield from SIMC  $\rightarrow$  cross section weighted, includes target thickness, RC, etc.

 $\sigma_{model} \rightarrow$  usually from stand-alone calculation. Difficult (although not impossible) to get cross section at particular point from simc. Crucial that stand-alone calculation matches model in simc exactly



#### **Extraction of cross sections - DIS**

$$\left[\frac{d\sigma(x_0, Q_0^2)}{dE_e d\Omega_e}\right]_{exp} = \frac{Y_{DATA}}{Y_{MC}(\sigma_{model})} \left[\frac{d\sigma(x_0, Q_0^2)}{dE_e d\Omega_e}\right]_{model}$$

 $Y_{data} \rightarrow$  efficiency-corrected # of counts normalized to charge

 $\rightarrow$  counts from root trees, charge and efficiencies from report file output from hcana

 $Y_{MC} \rightarrow$  More complicated than simc. Radiative corrections come from stand-alone program (outputs tables). MC yield from single-arm Monte-Carlo + cross section weights from tables. Must write program to combine RC weights + single-arm MC events

 $\sigma_{model} \rightarrow$  from same tables as RC (but using cross section at vertex instead)



## Analysis to-do list:

- Standard data analysis stuff
  - Reference time and time window cuts
  - Detector and beamline calibrations
  - Determine target boiling corrections, check other efficiencies
  - Determination of kinematic offsets
  - Determination of charge normalized, efficiency-corrected data yields

today

- Monte Carlo model iteration
- Extract cross sections/ratio  $\rightarrow$  Rosenbluth separations
- Publish
- High-priority for early in experiment running
  - Online reference time and time window cuts
  - Online calibrations
  - Event counter (simple, fast) to ensure we are taking adequate statistics
  - Normalization checks compare elastic and DIS yields to MC
  - Online target boiling checks
  - Optics checks (?)



## **Preparations for next meeting**

- Make sure you can access JLab computer systems remotely
  - Easiest to request access to linux VDI
  - https://jlab.servicenowservices.com/sp?sys\_kb\_id=dec16b0ddb7f0410ee4a3889fc961944&id =kb\_article\_view&sysparm\_rank=1&sysparm\_tsqueryId=db133dbc97e79a507d05bba6f053af cd
  - Submit request to <a href="mailto:helpdesk@jlab.org">helpdesk@jlab.org</a> if you don't already have access
- Get 2-factor token  $\rightarrow$  needed to access ifarm computers (another helpdesk request)
- Check if you are part of the c-rsidis group  $\rightarrow$  type "groups" on ifarm
  - If not, email Hanjie Liu (<u>hanjie@jlab.org</u>) to request to be added
  - This will give you access to the rsidis group and work disks
- Learn about loading modules:
  - <u>https://jlab.servicenowservices.com/scicomp?id=kb\_article\_view&sysparm\_article=KB001467</u>
     <u>1</u>
  - To access root on ifarm, you'll have to load the appropriate module
- If you like python, you can request access to the JLab instance of jupyterhub:
  - https://jupyterhub.jlab.org/
- RSIDIS elog: https://hallcweb.jlab.org/elogs/R-SIDIS+Experiment/

